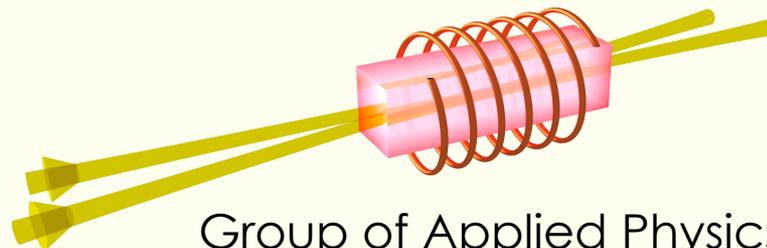


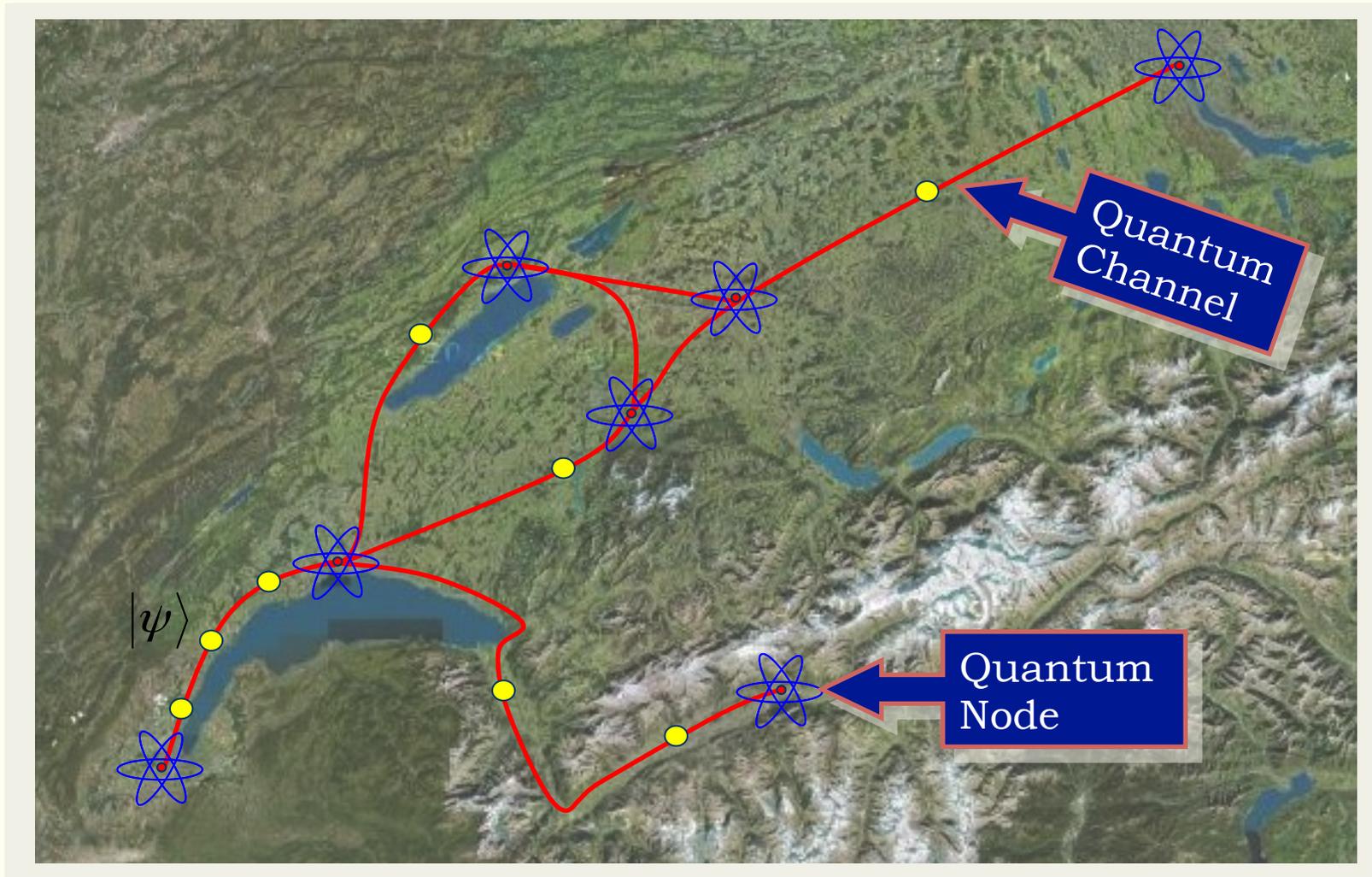
Generation of temporally multiplexed pairs of photons with controllable delay in a crystal

Cyril Laplane, Jean Etesse, Pierre Jobez,
Nicolas Gisin, Mikael Afzelius



Group of Applied Physics

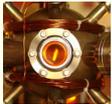
The quantum internet



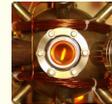
DLCZ-like elementary repeater link

How to distribute entanglement between two distant nodes ?

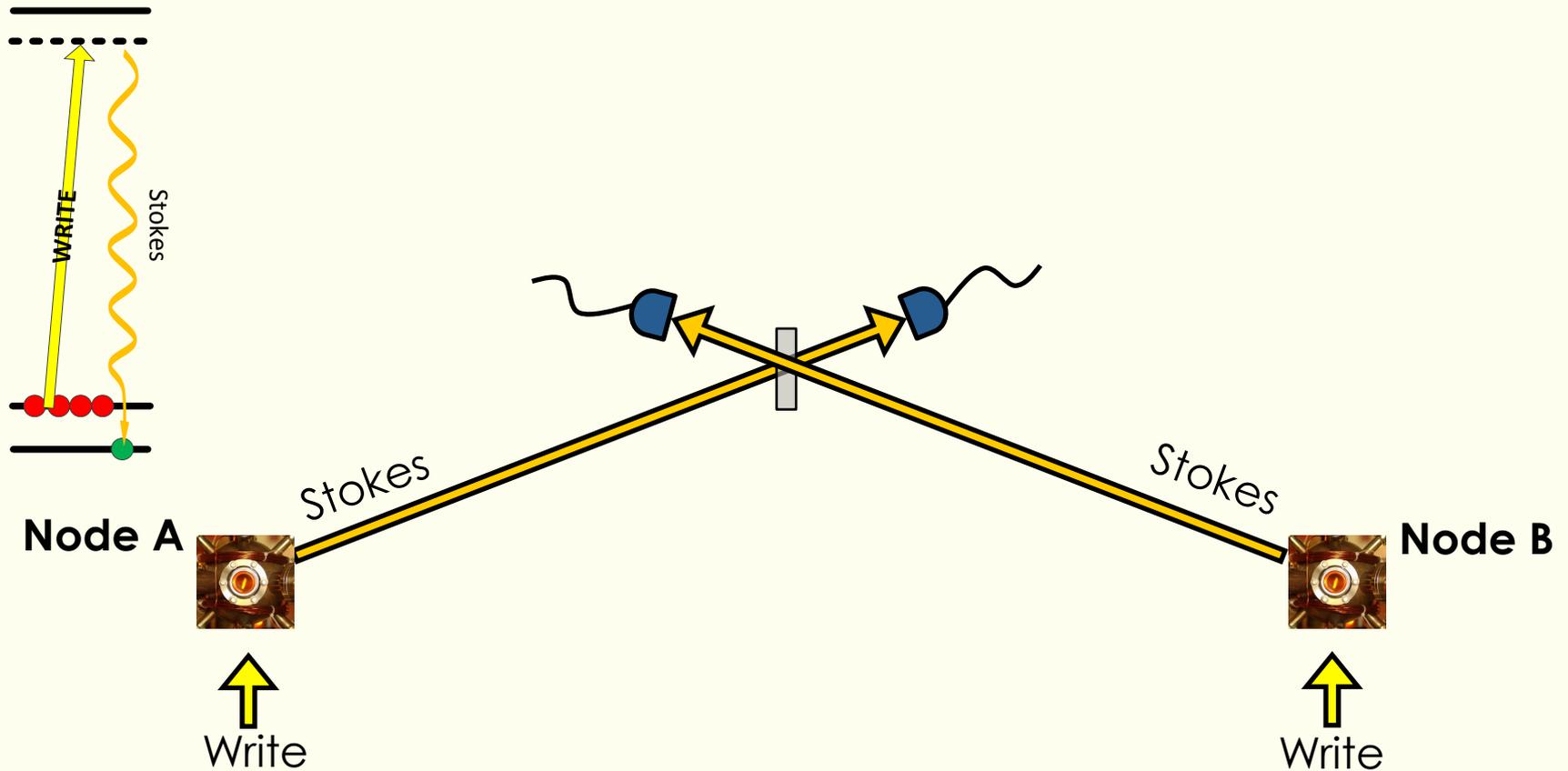
Node A



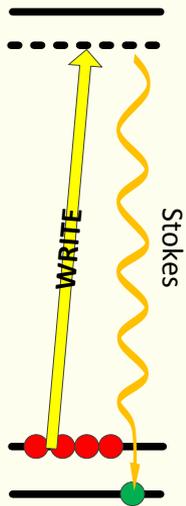
Node B



DLCZ-like elementary repeater link



DLCZ-like elementary repeater link

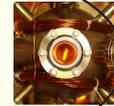


$$|\psi\rangle_{AB} = \frac{1}{\sqrt{2}} (|1\rangle_A |0\rangle_B + |0\rangle_A |1\rangle_B)$$

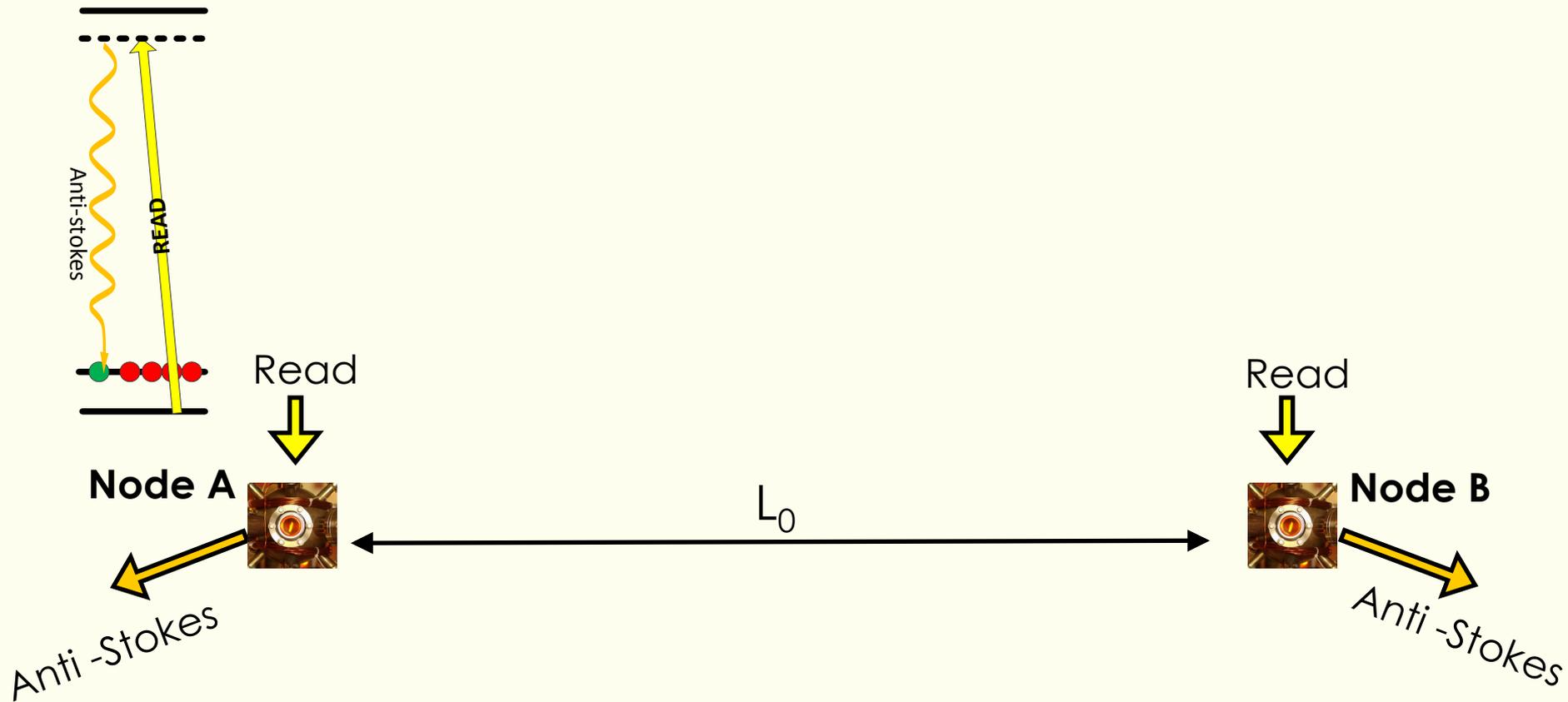
Node A



Node B



DLCZ-like elementary repeater link

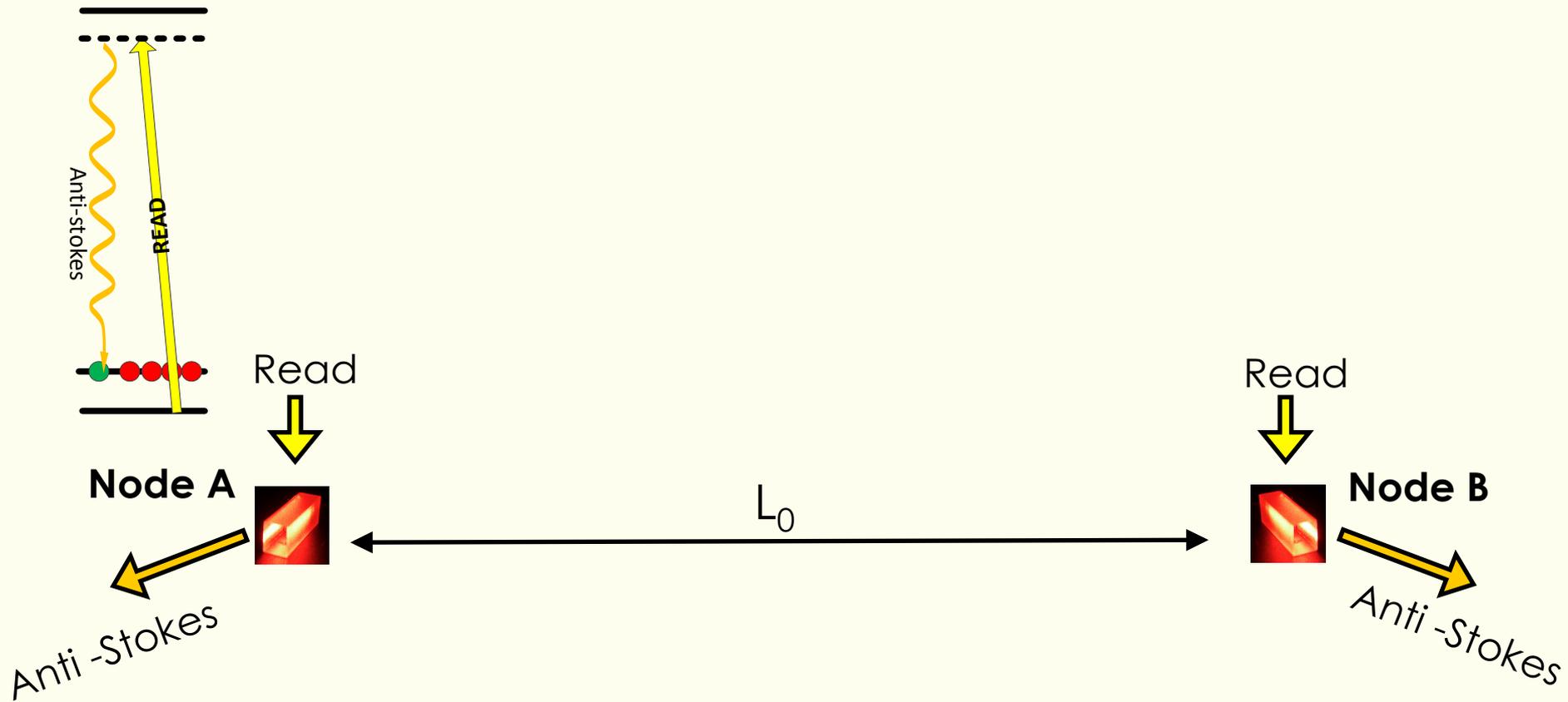


Single-mode storage: rate limited by c/L_0

N-mode storage: rate limited by $N.c/L_0$

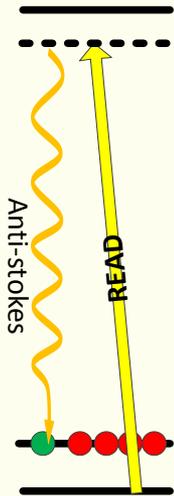
'Multimode DLCZ' scheme necessary for "fast" repeaters

DLCZ-like elementary repeater link



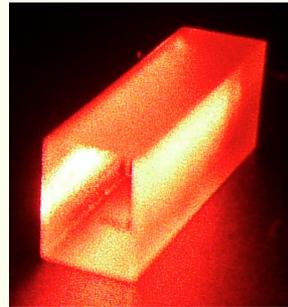
Our goal:
Solid state repeater link with REID crystals

DLCZ-like elementary repeater link



Our goal:
Solid state repeater link with REID crystals

Challenge:



Large inhomogeneous broadening

→ off-resonant fields for high cooperativity

But low electric dipole for REIDC ...

Outline

I – REID crystals as a quantum interface

II – The DLCZ protocol on REID crystals

III – Experimental results

Outline

I – REID crystals as a quantum interface

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Rare-earth ion doped crystals



lanthanum 57	cerium 58	praseodymium 59	neodymium 60	promethium 61	samarium 62	europium 63	gadolinium 64	terbium 65	dysprosium 66	holmium 67	erbium 68	thulium 69	ytterbium 70
La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb
138.91	140.12	140.91	144.24	[145]	150.36	151.96	157.25	158.93	162.50	164.93	167.26	168.93	173.04

λ (nm) = 606 880 580 1530 980

Properties of rare earths:

- Weak interaction with crystal environment
 - "atom" like energy structure for 4f-4f transitions
 - "frozen gas" of ions, no motional decoherence
- High number of stationary ions (10^7 - 10^{12})
strong light-matter coupling
- Long optical coherence times ($T < 4K$), $T_2^{\text{opt}} = 1 \mu\text{s} - 10 \text{ms}$
- Long spin coherence times ($T < 4K$), $T_2^{\text{hyp}} = 1 \text{ms} - 6 \text{h}^*$
- Large optical inhomogeneous broadenings 100 MHz – 10 GHz

*M. Zhong *et al.*, Nature **517**, 7533 177-180 (2015)

Rare-earth ion doped crystals



lanthanum 57 La 138.91	cerium 58 Ce 140.12	praseodymium 59 Pr 140.91	neodymium 60 Nd 144.24	promethium 61 Pm [145]	samarium 62 Sm 150.36	europium 63 Eu 151.96	gadolinium 64 Gd 157.25	terbium 65 Tb 158.93	dysprosium 66 Dy 162.50	holmium 67 Ho 164.93	erbium 68 Er 167.26	thulium 69 Tm 168.93	ytterbium 70 Yb 173.04
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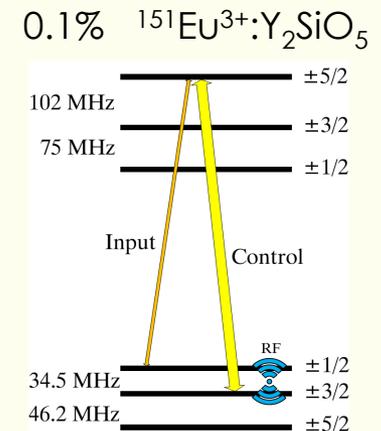
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P. Goldner



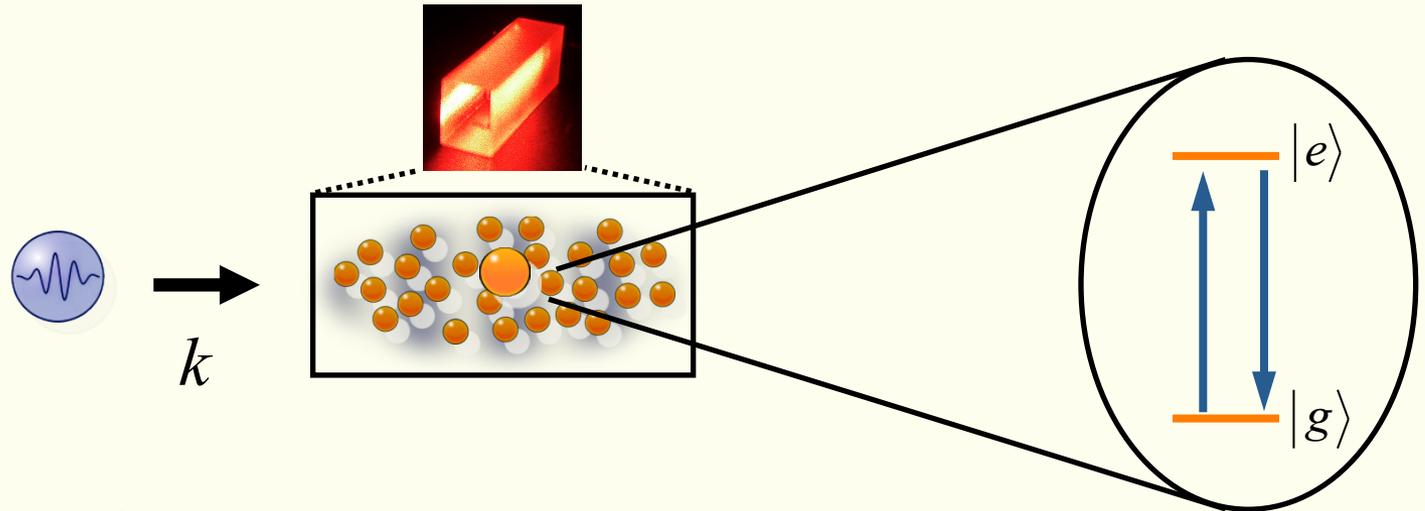
A. Ferrier



*M. Zhong et al., Nature **517**, 7533 177-180 (2015)

Inhomogeneity in REID crystals

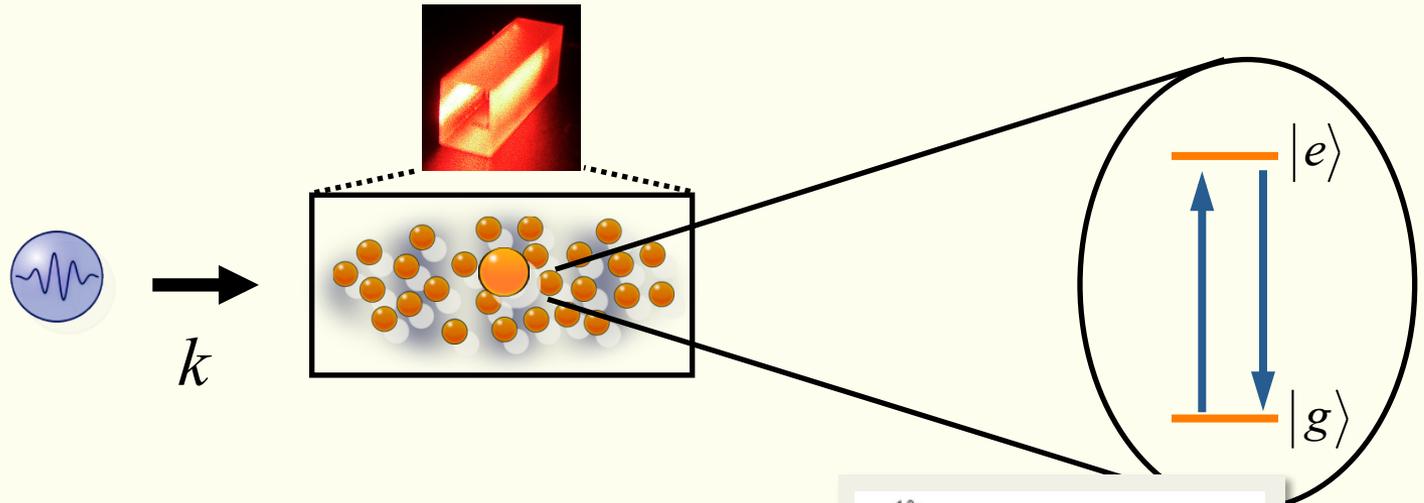
Absorption of an input photon by the ensemble creates a coherence...



$$|\psi\rangle \propto \sum_{i=1}^N c_i e^{i\delta_i t} e^{-ikz_i} |g_1 \dots e_i \dots g_N\rangle$$

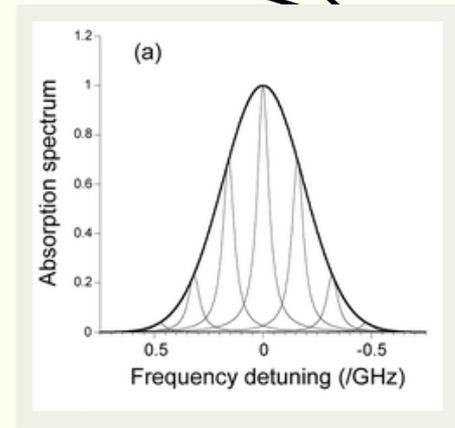
Inhomogeneity in REID crystals

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$$|\psi\rangle \propto \sum_{i=1}^N c_i e^{i\delta_i t} e^{-ikz_i} |g_1 \dots e_i \dots g_N\rangle$$

...which dephases quickly



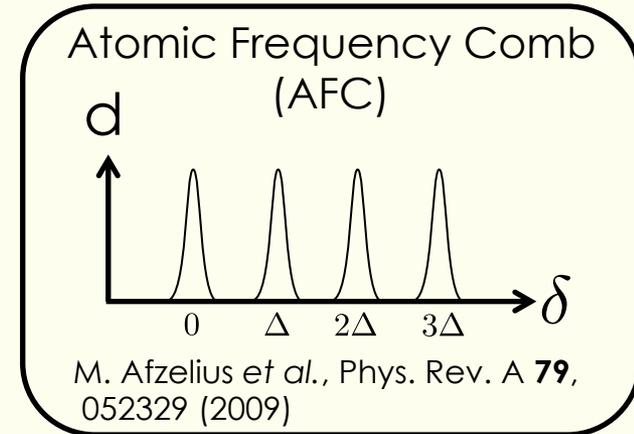
The AFC Protocol

The idea: only keep the atoms whose detuning is equal to

$$\delta_i = m_i \Delta$$

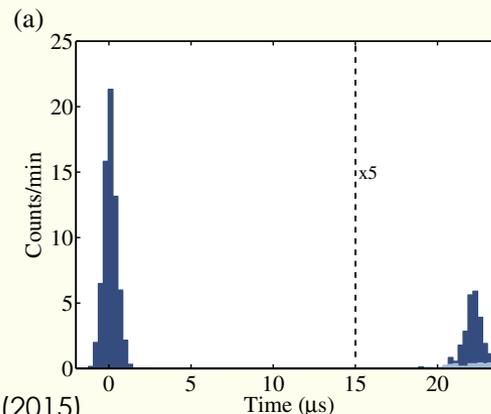
Such that the new state that evolves is

$$|\psi\rangle \propto \sum_{i=1}^{N'} c_i e^{im_i \Delta t} e^{-ikz_i} |g_1 \dots e_i \dots g_N\rangle$$



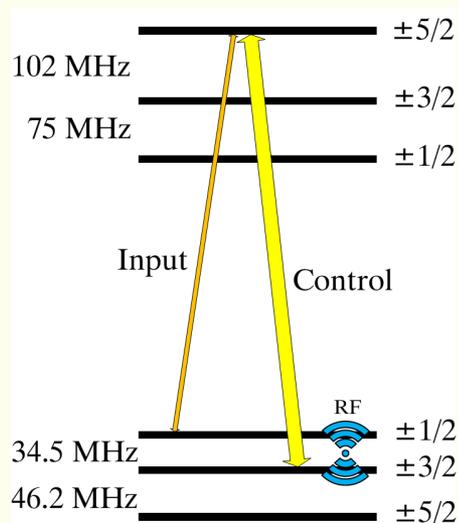
At time $t = \frac{2\pi}{\Delta}$, the selected atoms are back in phase !

Two level AFC
photon echo

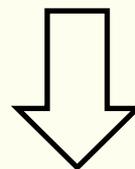


P. Jobez *et al.*, Phys. Rev. Lett. **114**, 230502 (2015)

The AFC Protocol



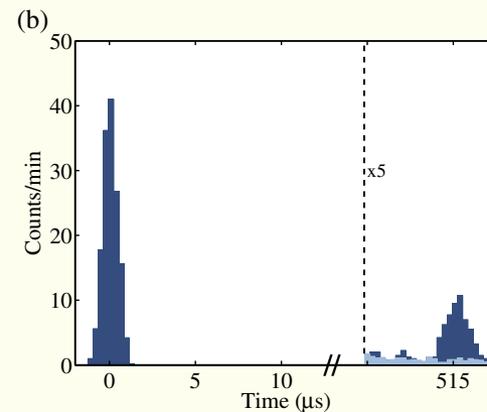
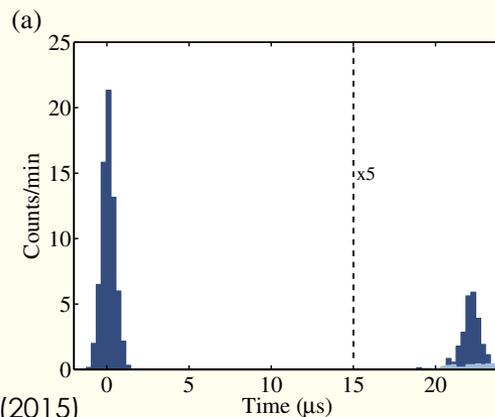
By adding a third level, the optical coherence can be transferred to a nuclear spin coherence



On-demand storage and retrieval of optical excitations

Three level AFC
photon echo

P. Jobez *et al.*, Phys. Rev. Lett. **114**, 230502 (2015)

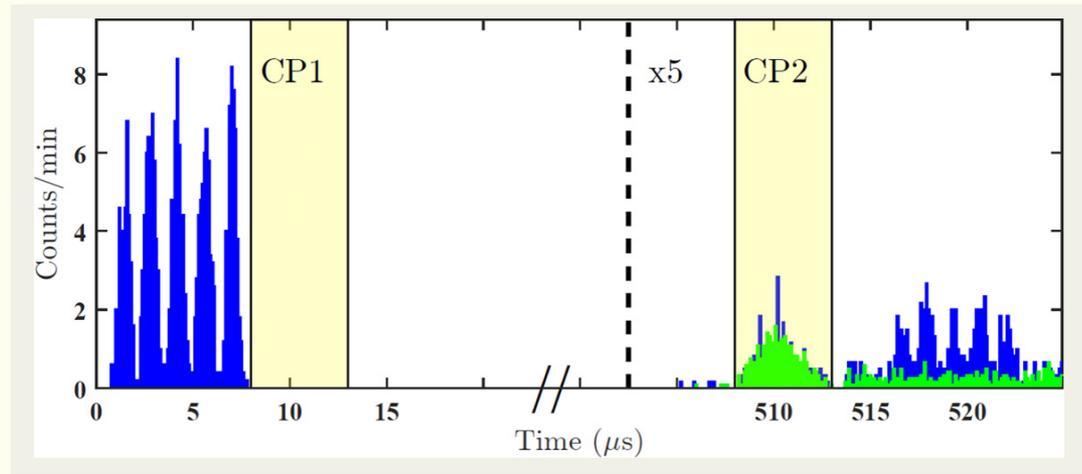


University of Geneva

Jean Etesse

The AFC Protocol

Force of the protocol: multimode capacity



Storage of five polarisation qubit modes

Outline

I – REID crystals as a quantum interface

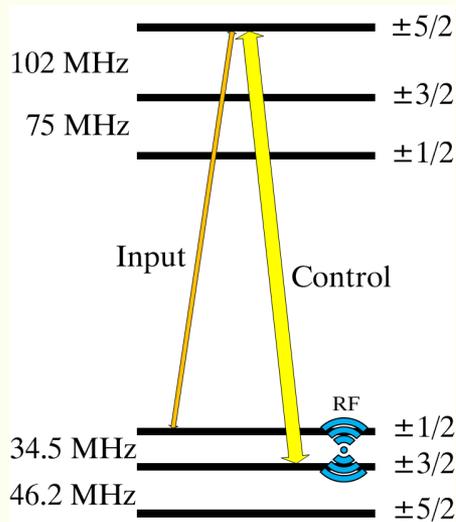
II – The DLCZ protocol on REID crystals

III – Experimental results

A homogeneous-like system

Thanks to the AFC rephasing effect, the system behaves like an homogeneous system at a precise instant, the instant of re-emission.

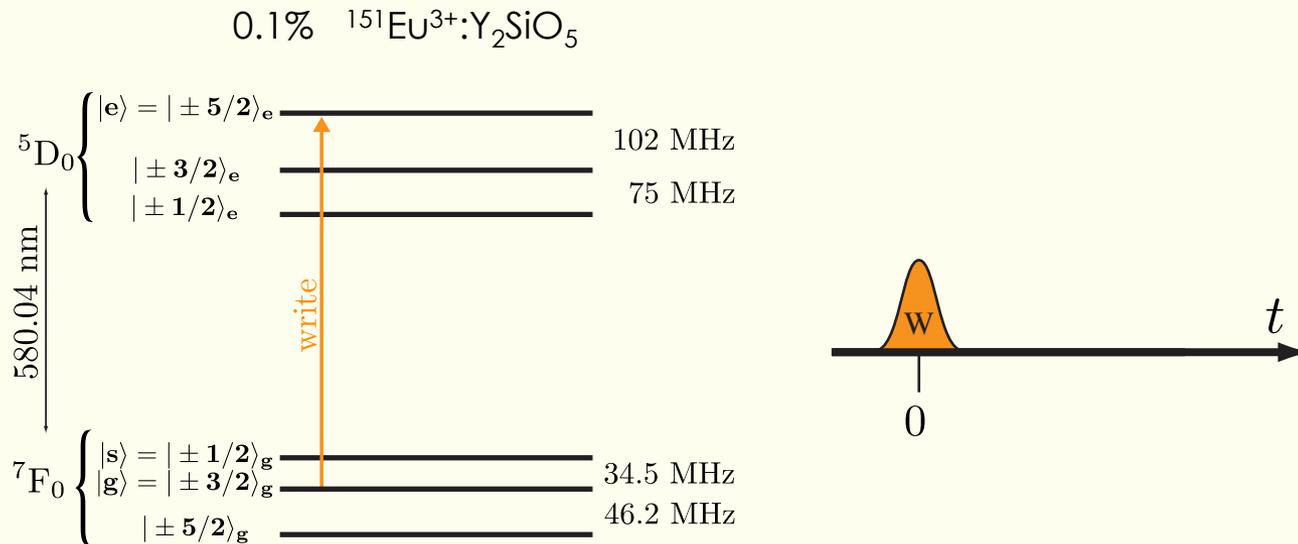
➔ Fields can be applied on resonance



Idea of the protocol: the control field is replaced by a spontaneous decay

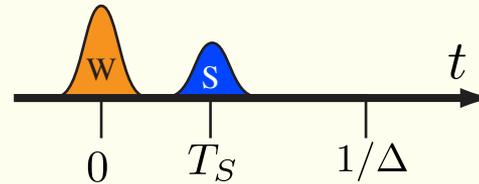
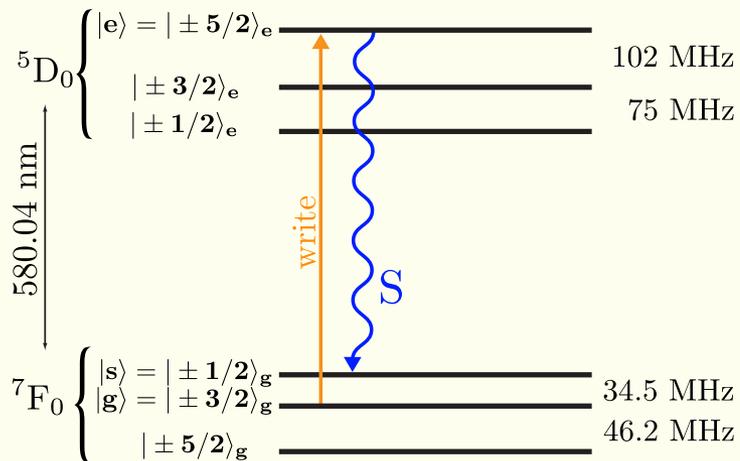
The AFC-DLCZ protocol

A weak 'write' pulse is sent on the **AFC** transition



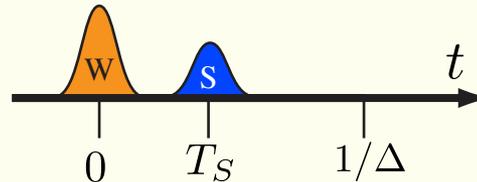
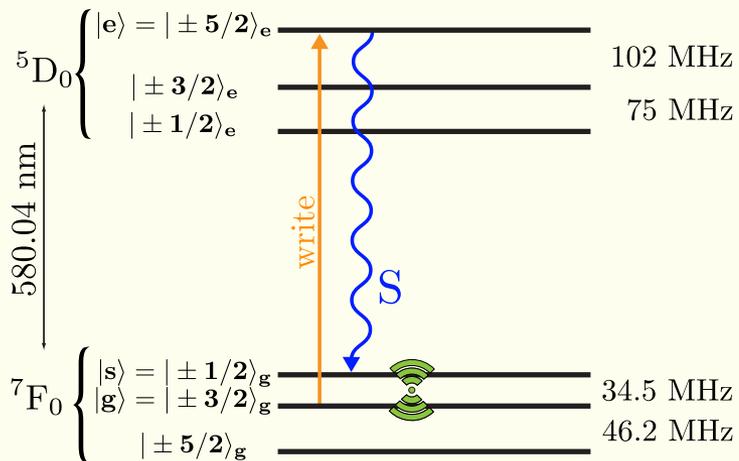
The AFC-DLCZ protocol

Stokes photons are spontaneously emitted on the strongest transition, before the AFC echo of the write pulse



The AFC-DLCZ protocol

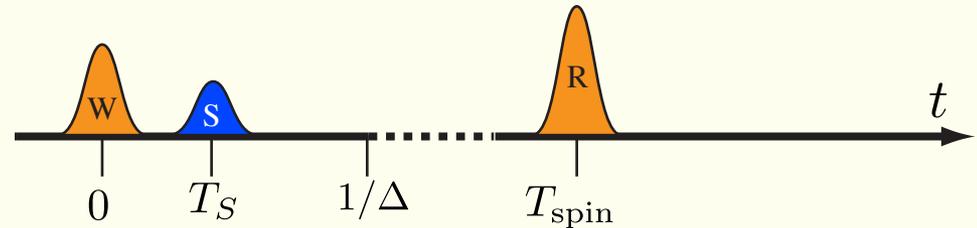
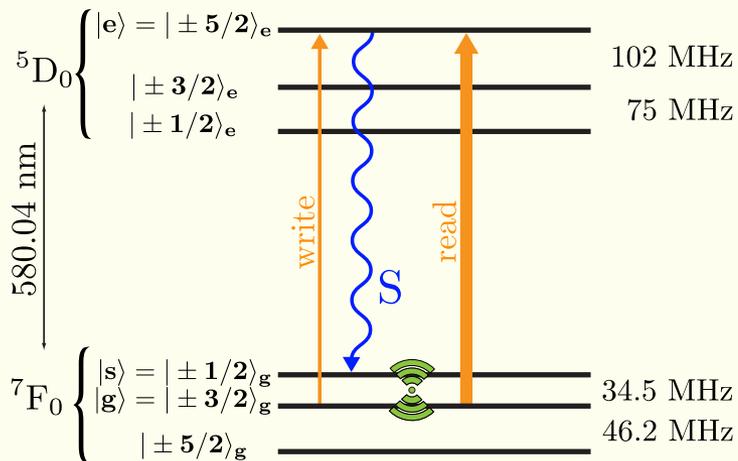
A radio-frequency pulse is sent to invert the populations on the $|s\rangle$ and $|g\rangle$ states...



The AFC-DLCZ protocol

A radio-frequency pulse is sent to invert the populations on the $|s\rangle$ and $|g\rangle$ states...

... so that read and write pulses have the same frequency...

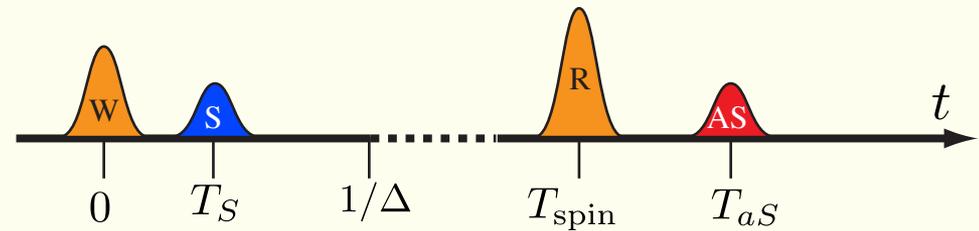
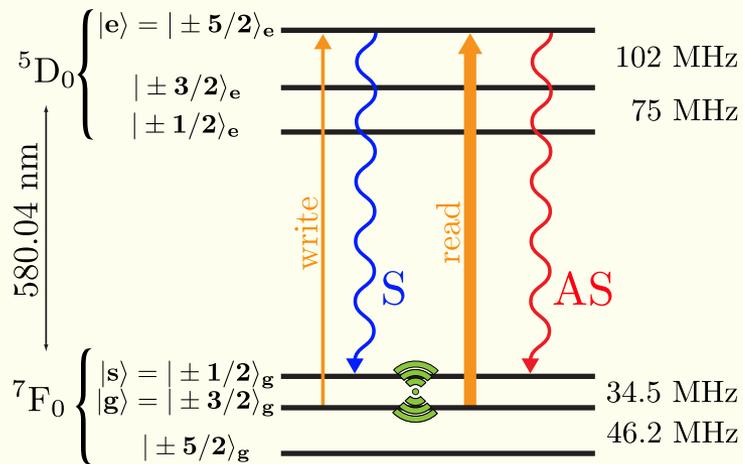


The AFC-DLCZ protocol

A radio-frequency pulse is sent to invert the populations on the $|s\rangle$ and $|g\rangle$ states...

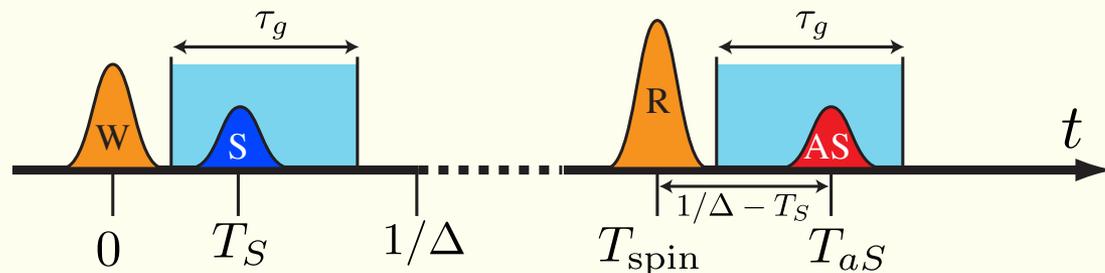
... so that read and write pulses have the same frequency...

... and Stokes and anti-Stokes as well.



The AFC-DLCZ protocol

The rephasing of the Stokes field occurs thanks to the AFC, leading to the emission of its echo: the anti-Stokes



Both are correlated in time: **temporal multiplexing capacity !**

Outline

I – REID crystals as a quantum interface

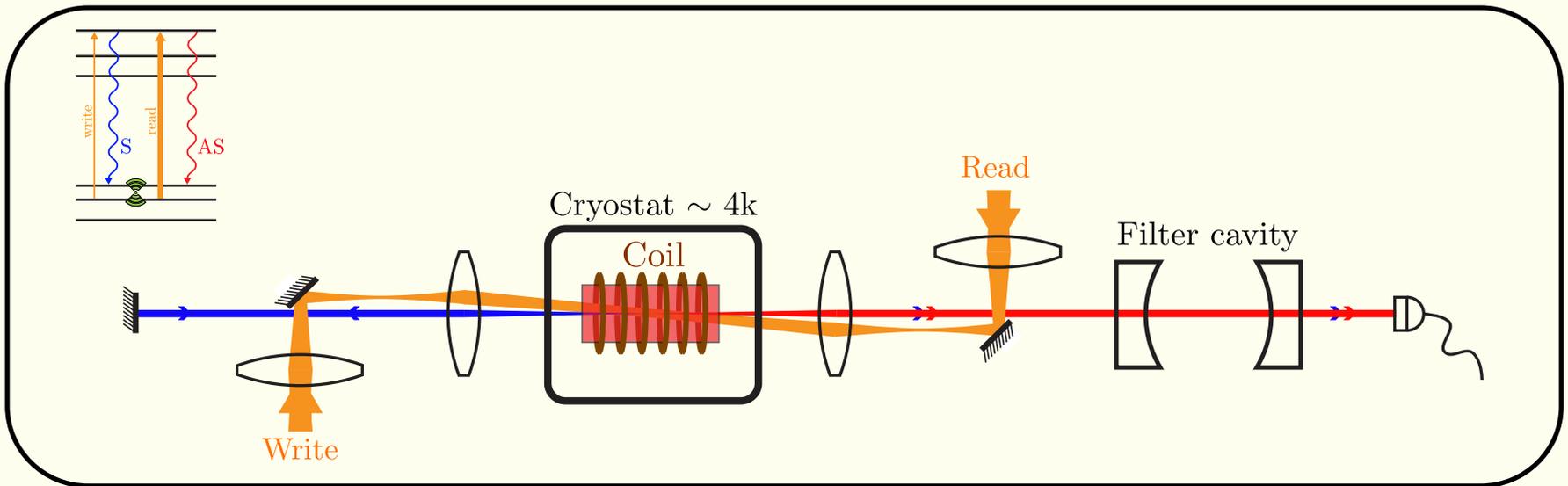
II – The DLCZ protocol on REID crystals

III – Experimental results

Experimental apparatus

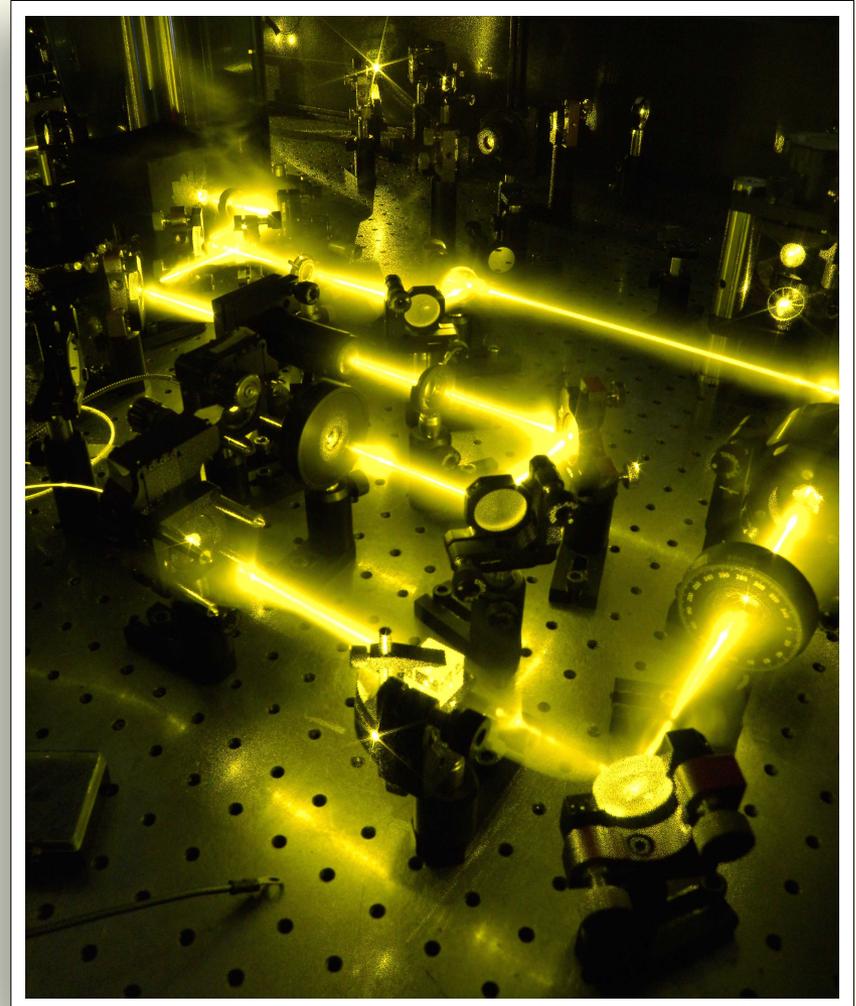
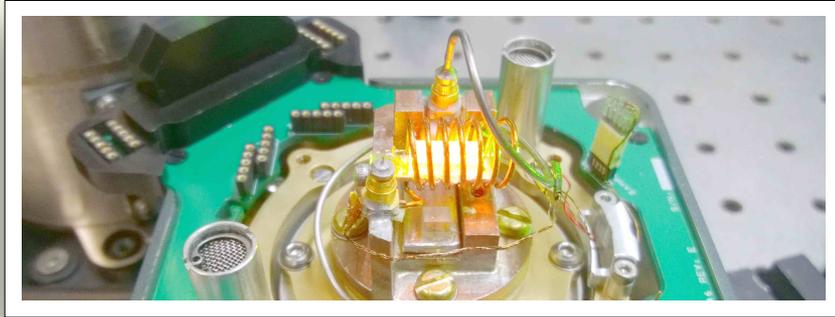
Need for a high-quality filtering:

- spatial angle between write/read and Stokes/aStokes
- Fabry-Perot cavity to detect only Stokes and aStokes

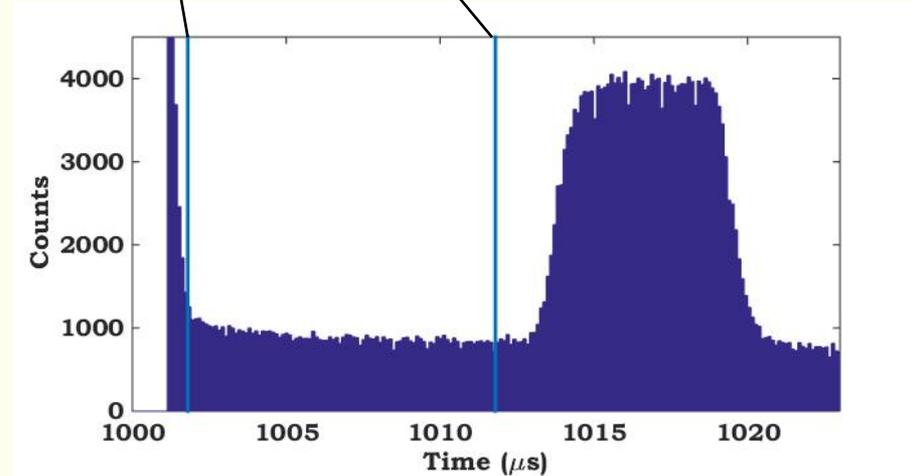
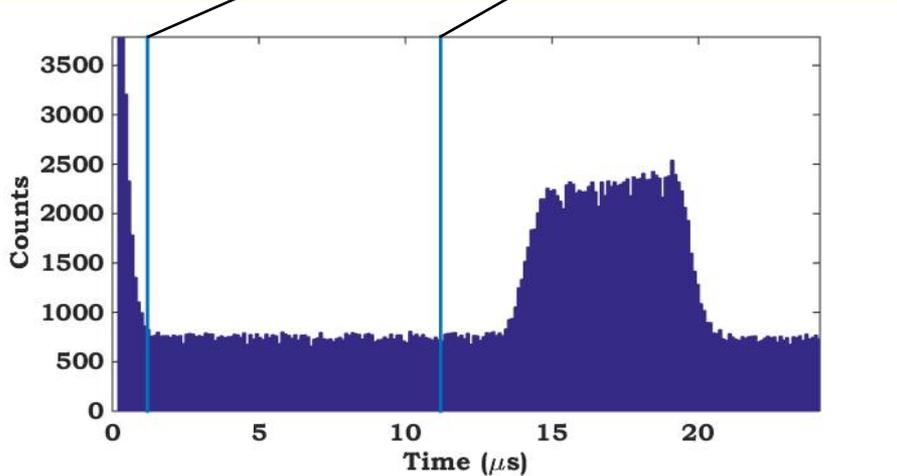
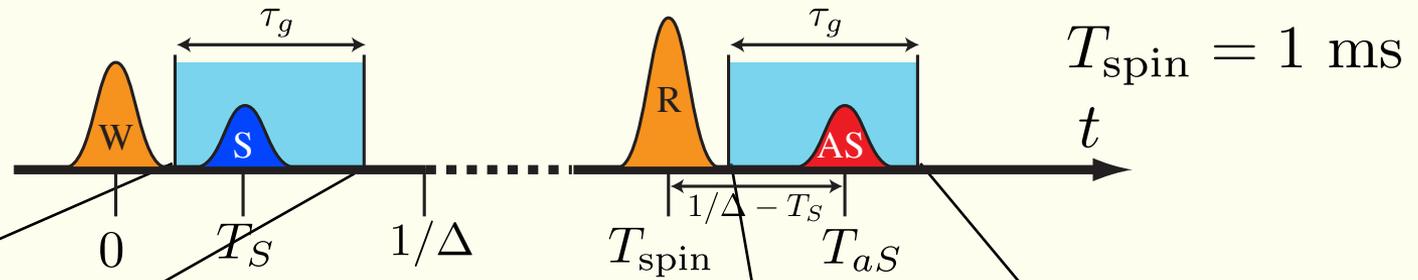


Double-pass configuration to double d and use only one cavity

Experimental apparatus



Results

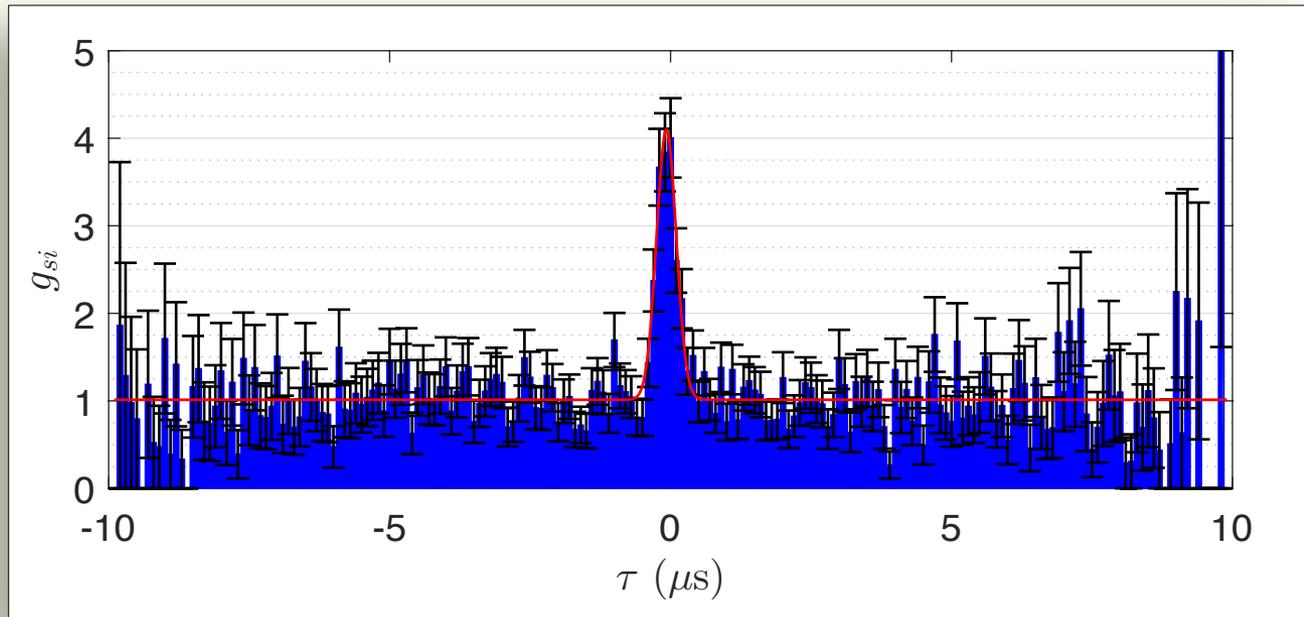


Integration time: 52.1069 h,
 2 626 190 experimental runs
 2 268 coincidences

Results

If we plot the cross-correlation function g_{si} as a function of

$$\tau = T_S + T_{aS} - T_{\text{spin}} - 1/\Delta$$

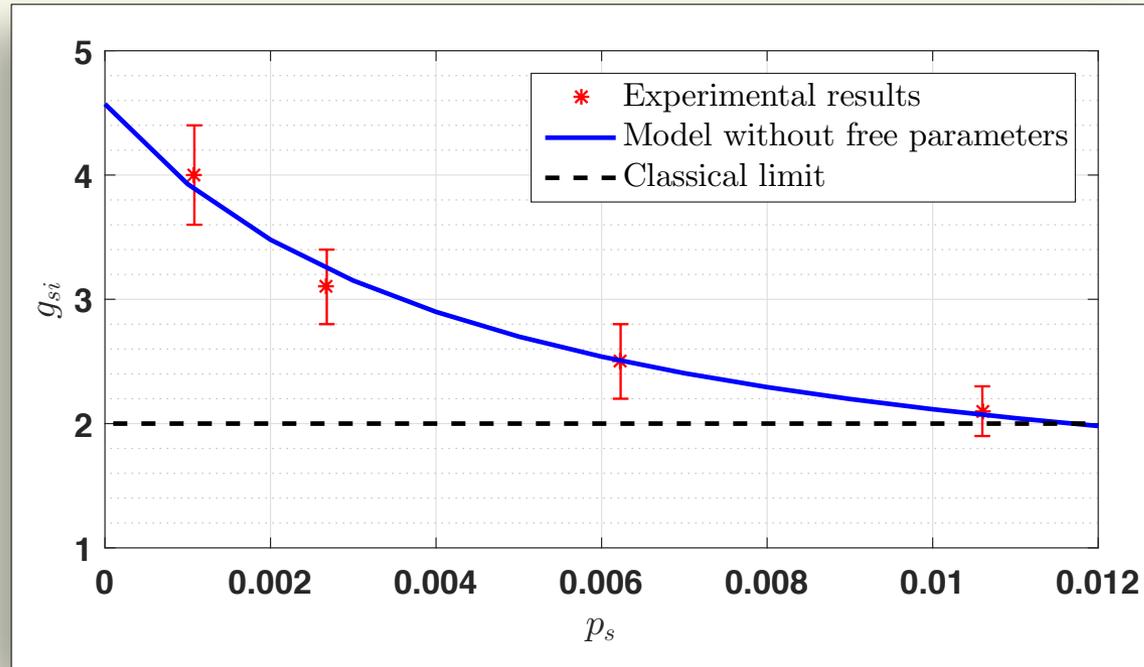


Fitted peak width: 414 ns FWHM

→ capacity of at least 10 modes in 10 μs gates

Results

Peak g_{si} as a function of the Stokes probability (100 ns bin)



$$\text{Model: } g_{si} = 1 + \frac{\eta}{(\eta + \beta)p_s + p_n}$$

Measured efficiency: $\eta = 0.5\%$

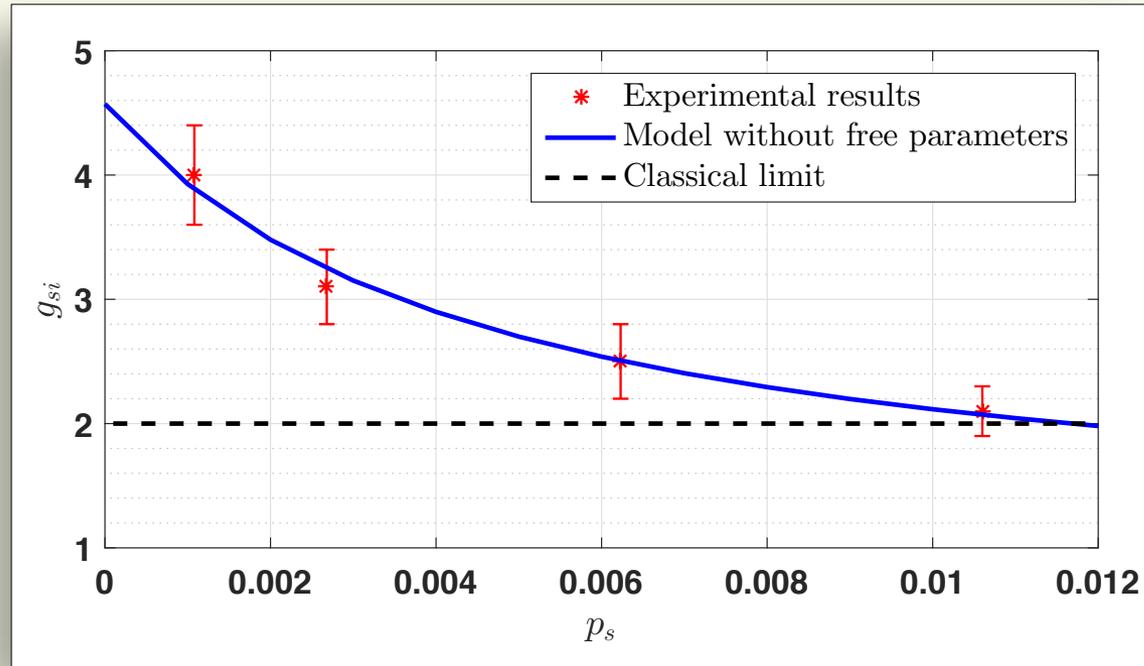
Intrinsic noise: $p_n = 0.14\%$

Noise dependant on p_s :

$$\beta = 0.3$$

Results

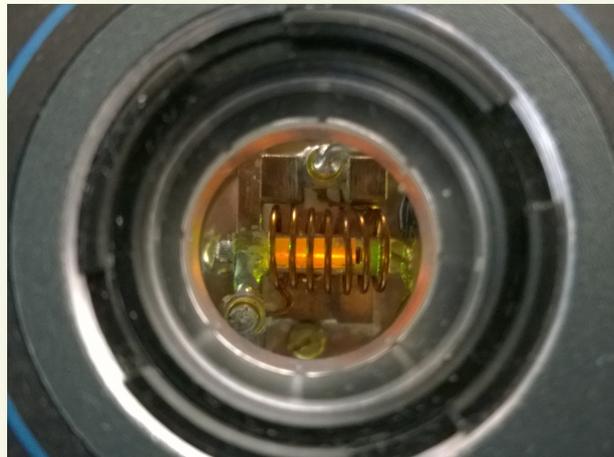
Peak g_{si} as a function of the Stokes probability (100 ns bin)



Non-classical correlation up to $p_s=10\%$ (in $1\mu s$)

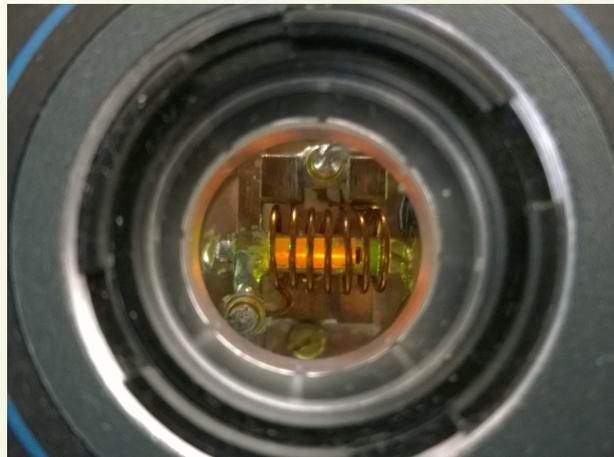
Conclusion

- An inhomogeneously broadened medium can be excited on resonance and still lead to a coherent emission
- Non-classical correlation between two single photons separated by 1 ms were demonstrated
- Temporal multiplexing up to 10 modes



Perspectives

- Measure the autocorrelation of Stokes and aStokes photons to prove the non-classical correlation un-ambiguously
- Explore new configurations to explain the low efficiency
- Prove the generation of delayed energy-time entanglement



Quantum memory group

