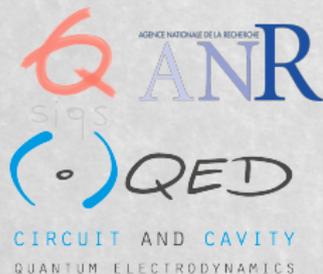


Long coherence times for Rydberg qubits on a superconducting atom chip

PHYSICAL REVIEW A **90**, 040502(R) (2014)

Thanh Long NGUYEN



QuantumGDR
Lyon 2014



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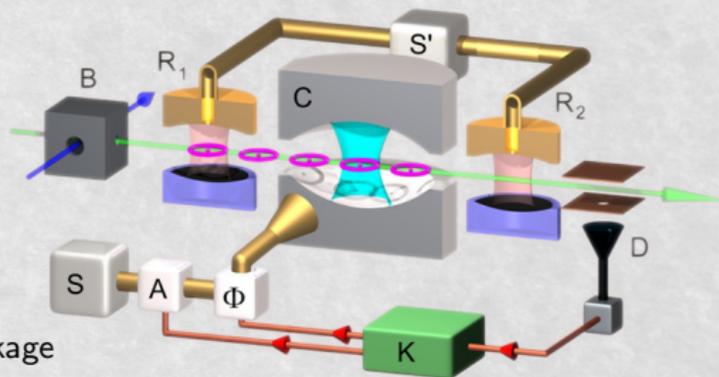
COLLÈGE
DE FRANCE
1530



INTRODUCTION

CQED EXPERIMENT

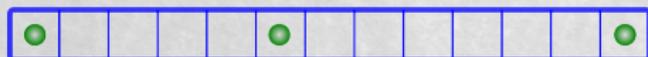
- QND measurement
- Fock state in a box
- Schrödinger cat



Hypothesis: single atoms on each package



Poisson distribution



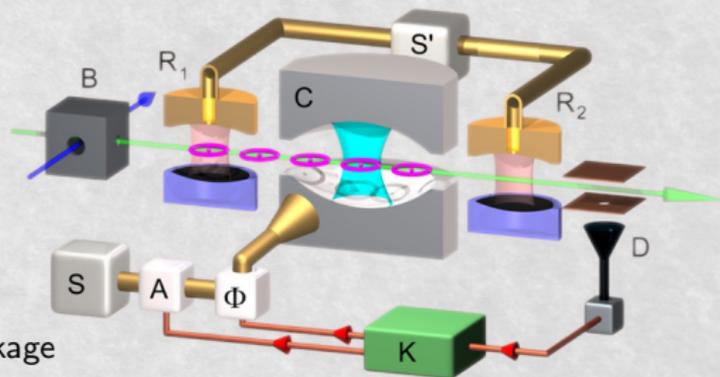
$\bar{n} \ll 1 \Rightarrow$ long data acquisition time

A deterministic single atom source

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16 JULY 2001

Dipole Blockade and Quantum Information Processing in Mesoscopic Atomic EnsemblesM. D. Lukin,¹ M. Fleischhauer,^{1,2} and R. Cote³¹*ITAMP, Harvard-Smithsonian Center for Astrophysics, Cambridge, Massachusetts 02138*²*Fachbereich Physik, Universität Kaiserslautern, D-67663 Kaiserslautern, Germany*³*Physics Department, University of Connecticut, Storrs, Connecticut 06269*

L. M. Duan, D. Jaksch, J. I. Cirac, and P. Zoller

Institut für Theoretische Physik, Universität Innsbruck, A-6020 Innsbruck, Austria

(Received 7 November 2000; published 26 June 2001)

We describe a technique for manipulating quantum information stored in collective states of mesoscopic ensembles. Quantum processing is accomplished by optical excitation into states with strong dipole-dipole interactions. The resulting "dipole blockade" can be used to inhibit transitions into all but singly excited collective states. This can be employed for a controlled generation of collective atomic spin states as well as nonclassical photonic states and for scalable quantum logic gates. An example involving a cold Rydberg gas is analyzed.

DOI: 10.1103/PhysRevLett.87.037901

PACS numbers: 03.67.Lx, 03.75.Fi, 42.50.Gy, 73.23.-b

Recent advances in quantum information science have opened a door for a number of fascinating potential applications ranging from the factorization of large numbers and secure communication to nanosensing techniques with

optical wavelengths. Combined with the exceptional degree of control that is typical for quantum optical systems and long coherence times, this allows one to considerably alleviate many stringent requirements for

Blockade of Rydberg excitation on a superconducting atom chip

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M. Saffman and T. G. Walker

Department of Physics, University of Wisconsin, 1150 University Avenue, Madison, Wisconsin 53706

(Received 14 March 2002; published 16 December 2002)

We discuss the application of dipole blockade techniques for the preparation of single-atom and single-photon sources. A deterministic protocol is given for loading a single atom in an optical trap, as well as ejecting a controlled number of atoms in a desired direction. A single-photon source with an optically controlled beamlike emission pattern is described.

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PACS number(s): 32.80.Qk, 03.67.Dd

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Quantum-information science builds on the laws of quantum mechanics to transmit, store, and process information in varied and powerful ways. Advances in this field rely on our ability to manipulate coherently isolated quantum objects while eliminating incoherent interactions with the surrounding environment. It was proposed two decades ago that the

in the rotating wave approximation, and $V_{dd}/\hbar = \sum_{j,k>j} \Delta_{jk} |r_j r_k\rangle \langle r_j r_k|$ describes the dipole-dipole interaction of two excited atoms. Here $\Omega_j = -\langle r_j | \hat{d} | b_j \rangle \mathcal{E}(\mathbf{r}_j) / \hbar$, where \hat{d} is the dipole moment operator, and the position dependent optical field is $E(\mathbf{r}_j) = [\mathcal{E}(\mathbf{r}_j)/2] e^{-i\omega t} + c.c.$ The dipole-dipole shift in the case of dipole moments aligned

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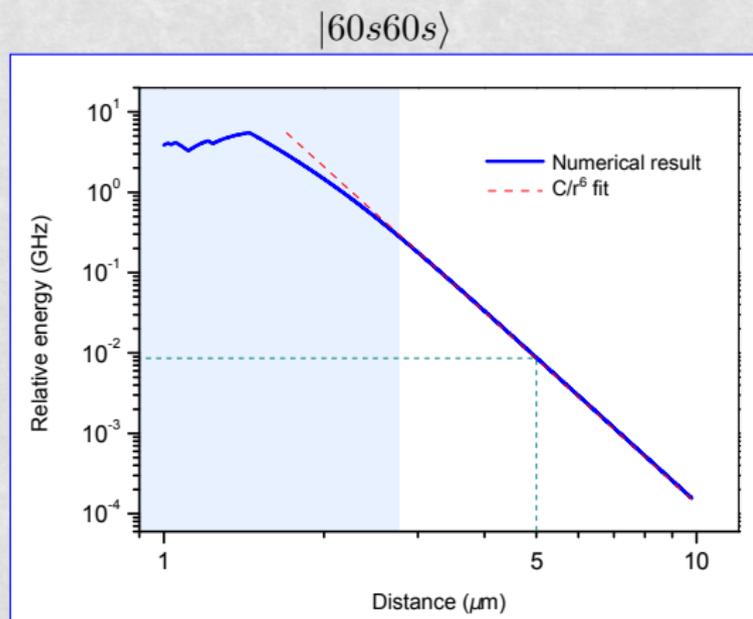
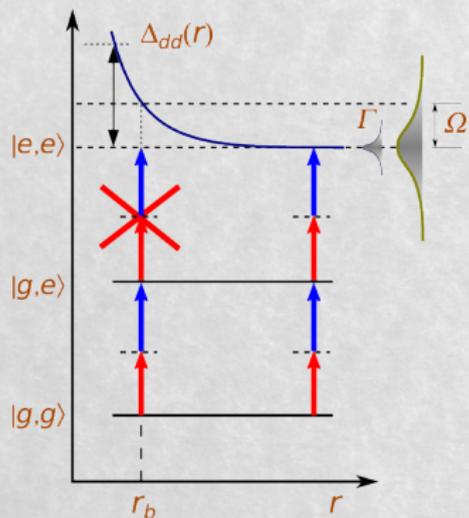
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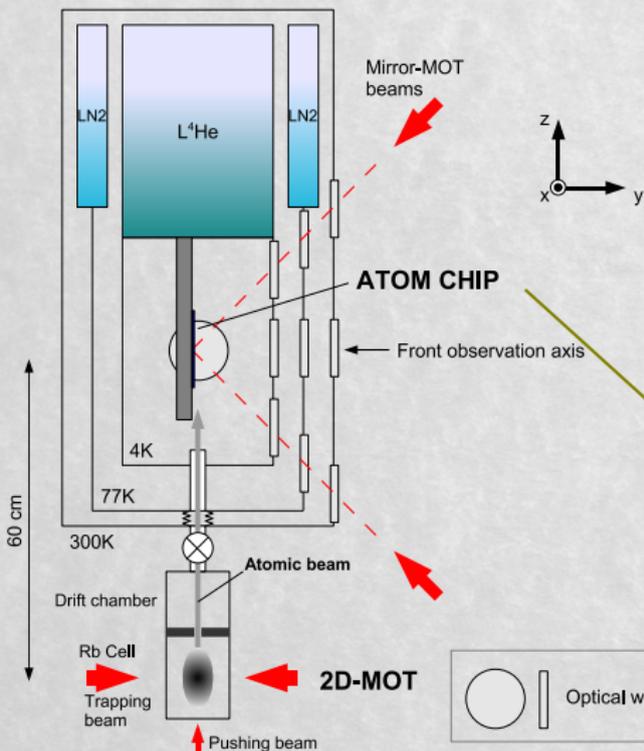
Taming electric field near a superconducting atom chip

DIPOLE BLOCKADE



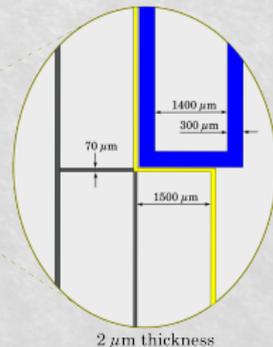
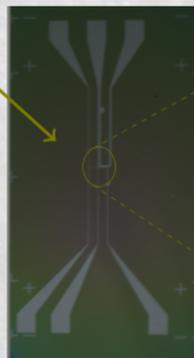
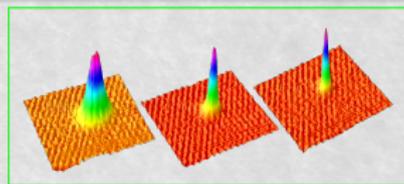
- Rydberg atoms: $n \gg 1 \Rightarrow$ huge electric dipole matrix element $d \propto n^2$
- Dipole interaction \Rightarrow shift of energy levels
- $\Delta_{dip} >$ excitation line-width \Rightarrow blockade regime
- Radius of blockade r_b

EXPERIMENT SETUP

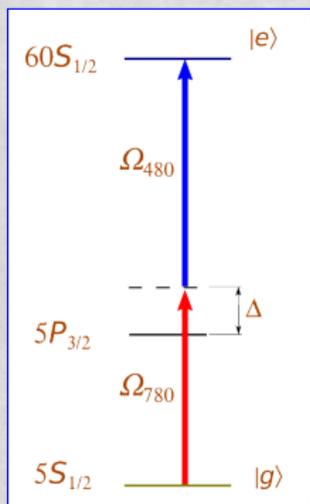


SUPERCONDUCTING ATOM CHIP

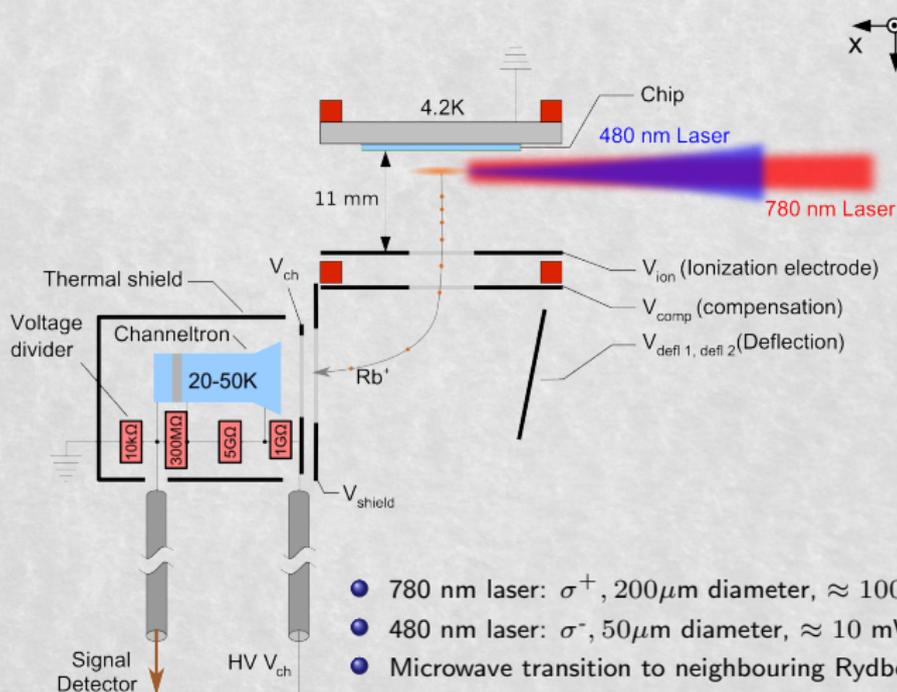
- Nb wires: superconducting
- Gold coated
- On chip 3D MOT & IP magnetic trap
- Varied position and temperature



RYDBERG EXCITATION & DETECTION



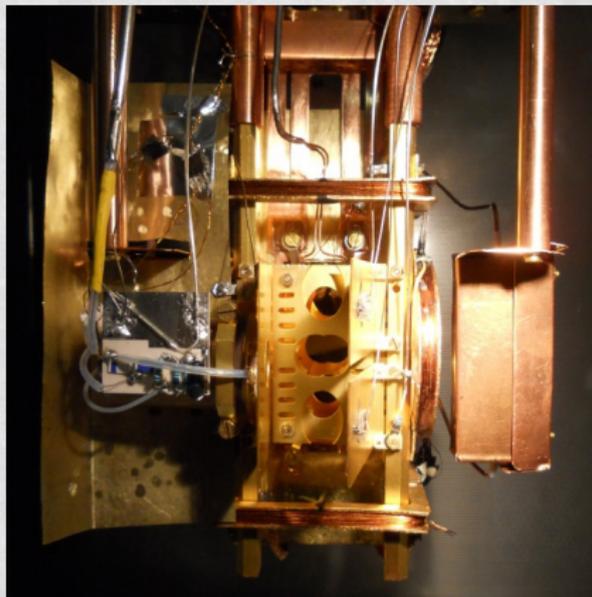
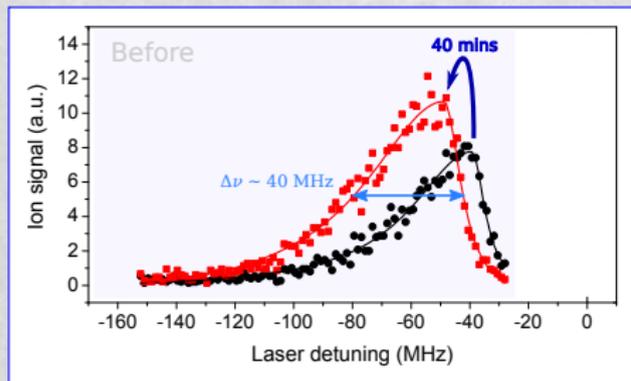
$$\Delta = 540\text{MHz}$$



- 780 nm laser: σ^+ , $200\mu\text{m}$ diameter, $\approx 100\text{ nW}$
- 480 nm laser: σ^- , $50\mu\text{m}$ diameter, $\approx 10\text{ mW}$
- Microwave transition to neighbouring Rydberg levels
- State selective detection
- Compensation for perpendicular residual electric field

STARK EFFECT AND OPTICAL TRANSITION

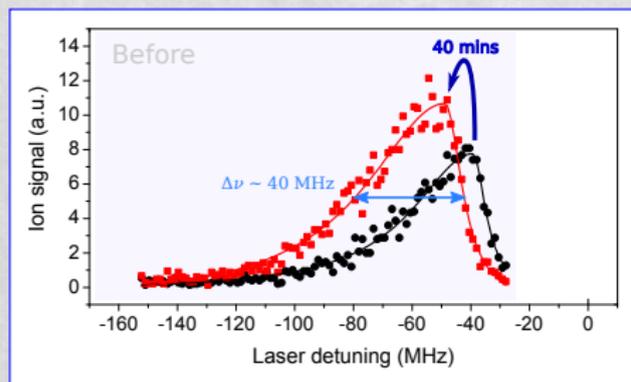
Spatial and temporal fluctuation of stray field due to **accumulation of charges** on chip surface



J. D. Carter, O. Cherry and J. D. D. Martin, PRA **86**, 053401 (2012): gradient of 100 (V/cm²) at 500 μm from the chip surface + temporal drifts
 A. Tauschinsky et al., PRA **81**, 063411 (2010): field of 1 V/cm at 200 μm, temporal drift due to slow Rb deposition

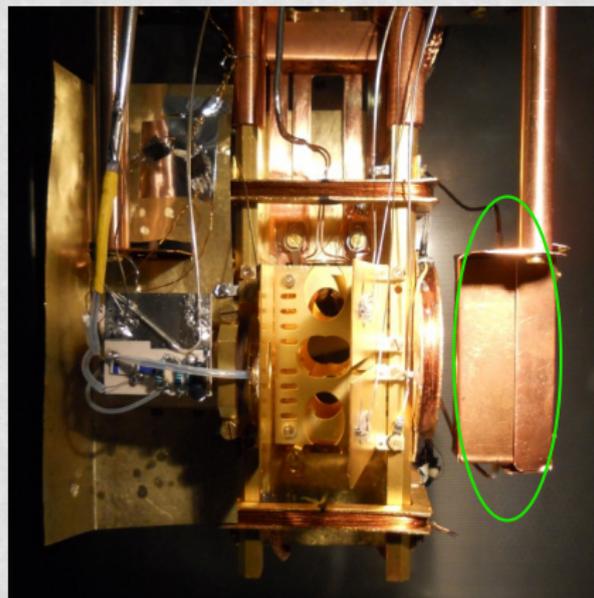
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SOLUTION IN OUR CASE

Installation of **Rb dispensers** inside the **cryogenic environment** for Rb deposition on a large surface of the chip → **saturation** of the slow **Rb deposit** (~ 90 nm layer)

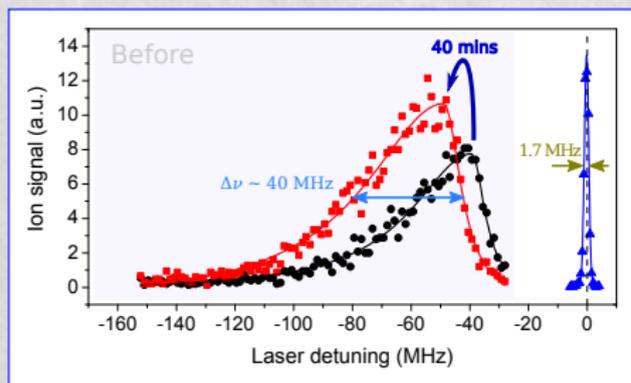


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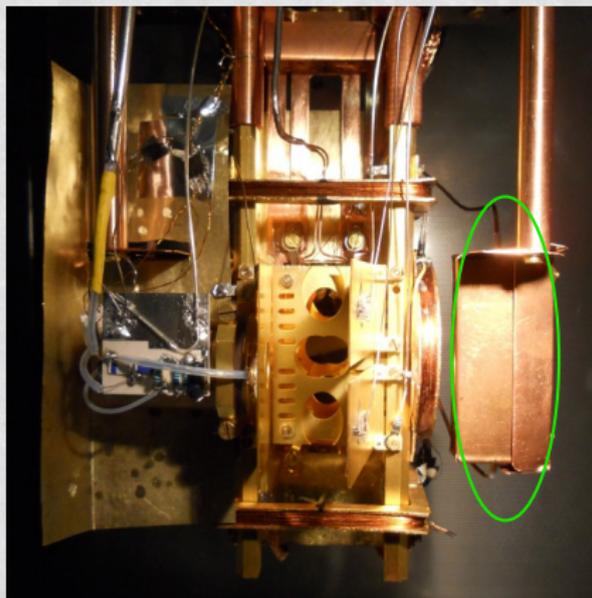
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ELECTRIC FIELD MEASUREMENT USING MICROWAVE

Precise knowledge of Stark shift allows to measure the residual electric field

$$\hat{H}\psi = \hat{H}_0 + e\vec{F} \cdot \vec{r}$$

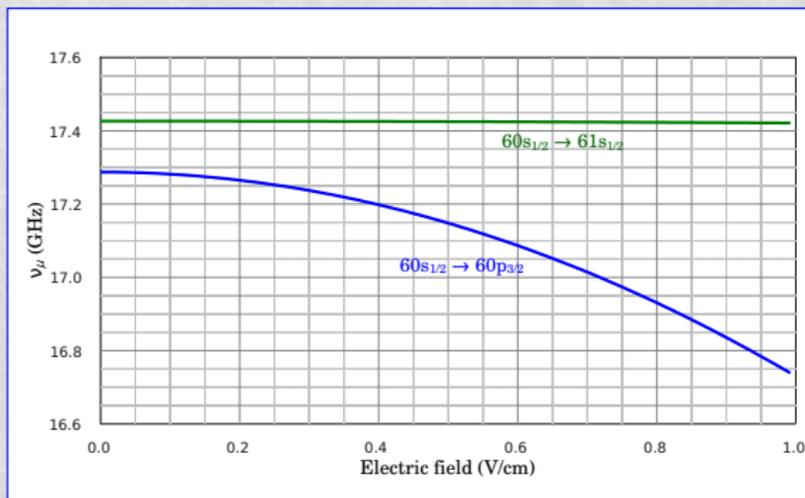
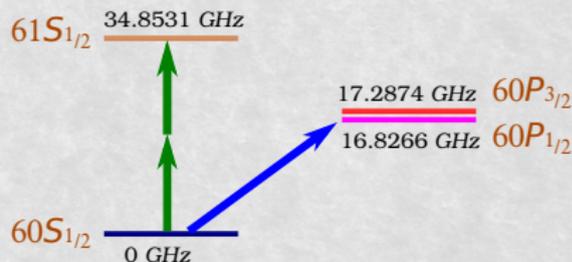
$$\Delta\nu = A|\vec{F}|^2, A \propto n^7$$

$$A_{60s1/2} = -92.2 \text{ MHz}$$

$$A_{61s1/2} = -102.8 \text{ MHz}$$

$$A_{60p3/2} = -649.6 \text{ MHz}$$

Microwave spectroscopy allows accurate probe of energy shifts
 → precise tool to measure electric field



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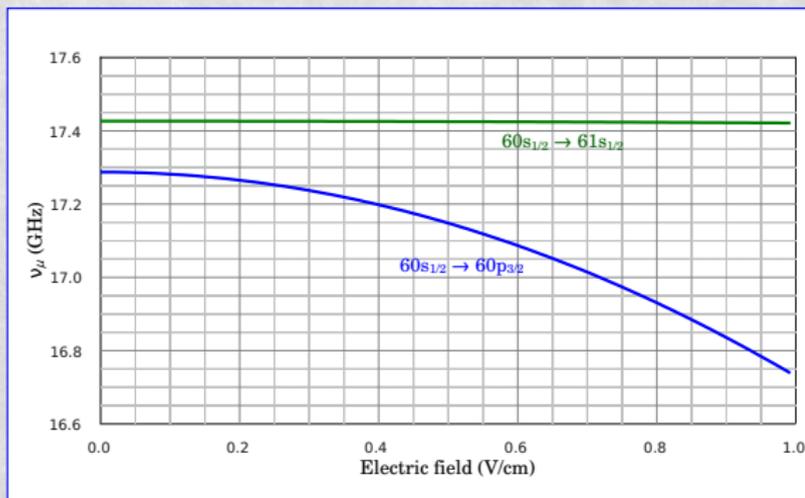
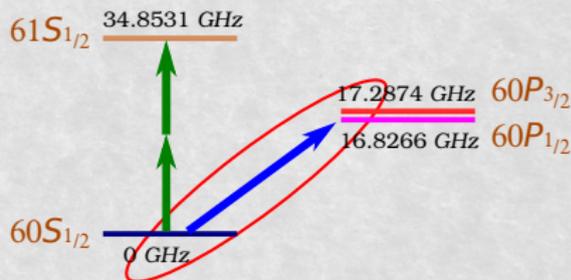
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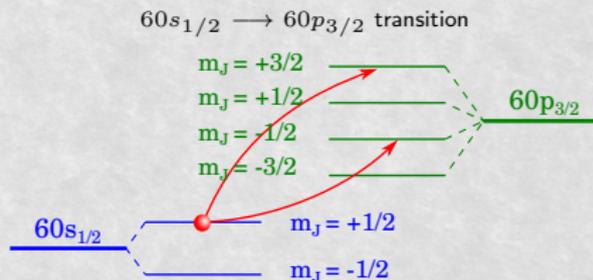
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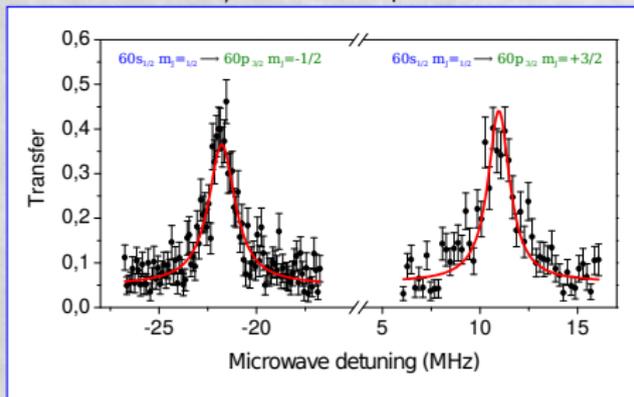


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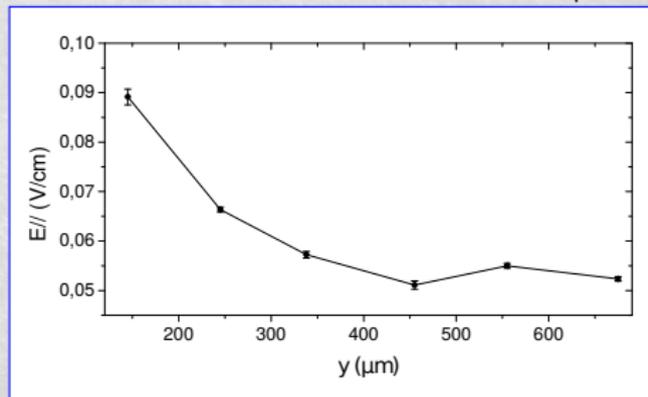
- Thermal cloud in a magnetic trap
- Taking into account Zeeman splitting \rightarrow We can deduce the electric field with certain precision



@455 μm from the chip surface

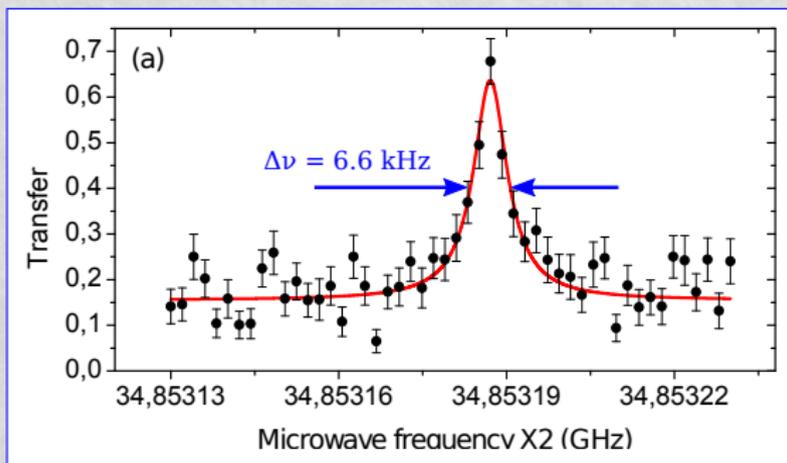
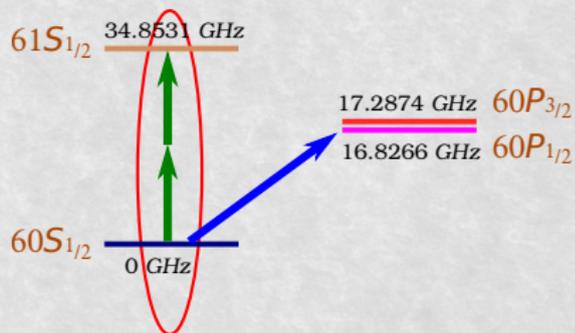


Residual electric field as function of distance to chip

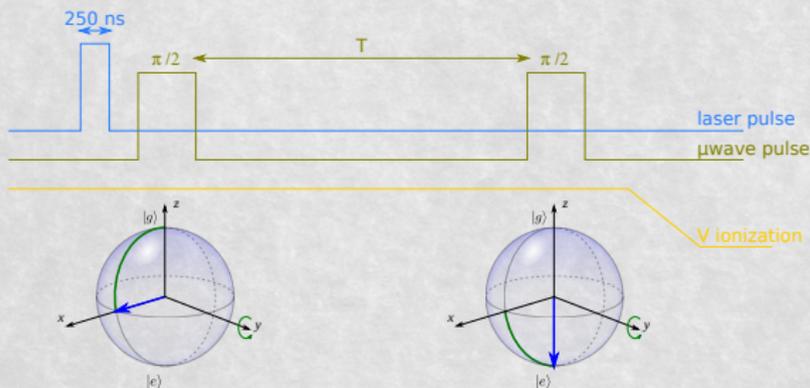


COHERENCE - SPECTRUM

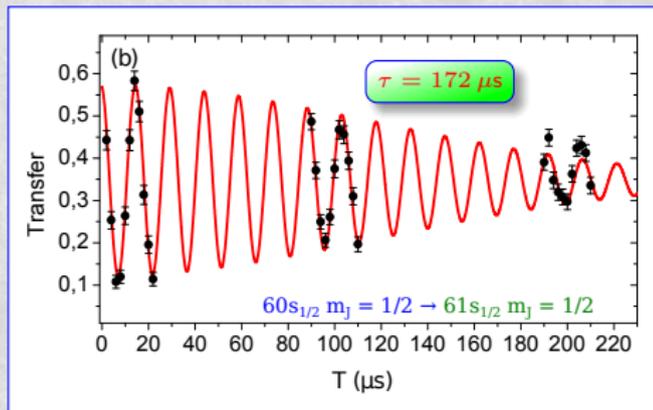
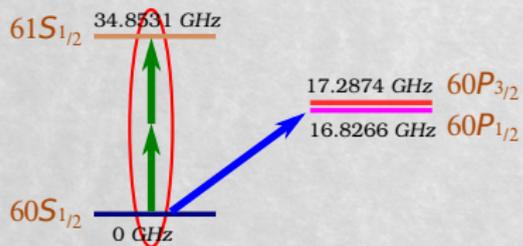
- Thermal cloud in a magnetic trap
 @450 μm from the chip
 $T \sim 1 \mu\text{K}$
 $\sim 60 \mu\text{m} \times 30 \mu\text{m} \times 30 \mu\text{m}$



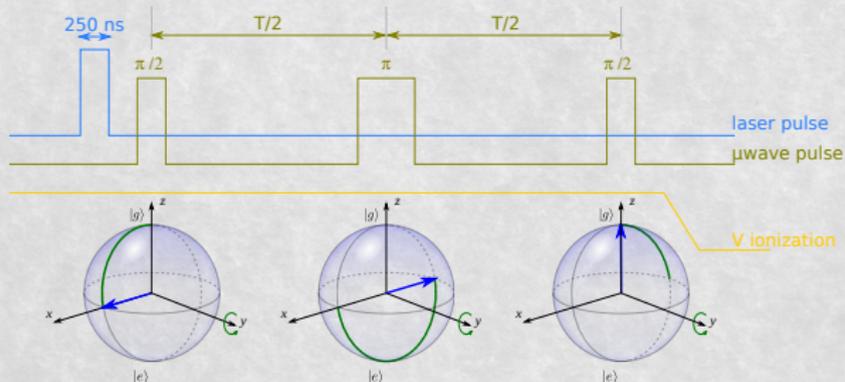
COHERENCE-RAMSEY INTERFERENCE



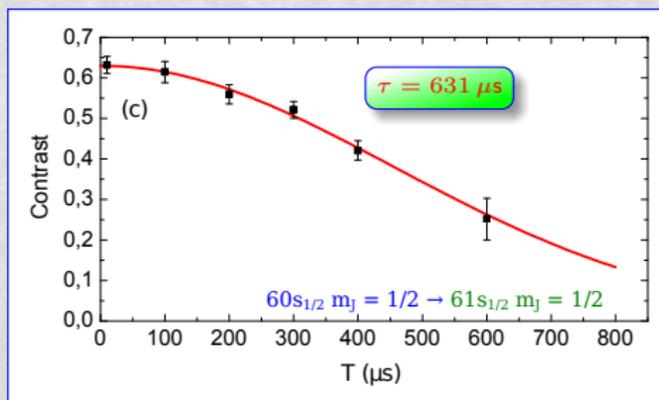
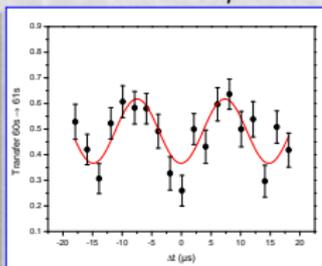
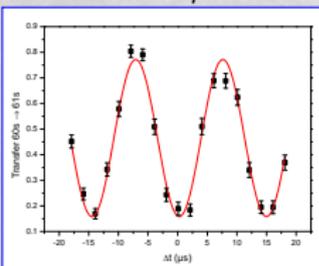
Thermal cloud @450 μm from the chip
 $\Delta\nu = 70$ kHz



COHERENCE-SPIN ECHO



Thermal cloud @450 μm from the chip
 Rydberg lifetime $\sim 210 \mu\text{s}$
 $\Delta\nu = 2 \times 35 \text{ kHz}$
 $T = 100 \mu\text{s}$ $T = 600 \mu\text{s}$



CONCLUSION

- **Rb dispenser** inside the cryogenic environment allowed a good **control of residual electric field** close to the chip surface
- Residual electric field were probed from **microwave spectroscopy** of nearby Rydberg transition
- **Coherence** times bigger than the Rydberg lifetime were demonstrated

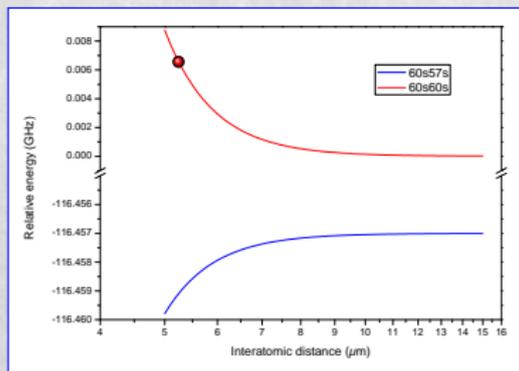
A step forward deterministic excitation of single Rydberg atoms

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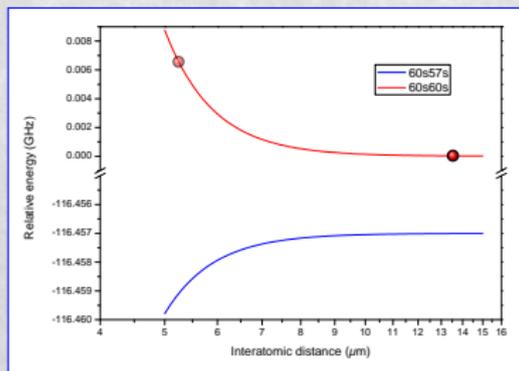
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EXPANSION OF A CLOUD OF RYDBERG ATOMS



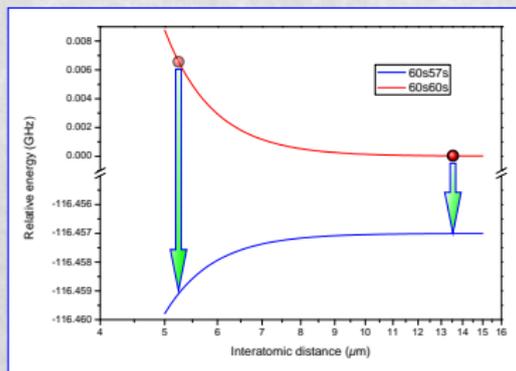
- Detune laser to give initial potential
- Use microwave to map the spatial distribution of the Rydberg cloud

EXPANSION OF A CLOUD OF RYDBERG ATOMS

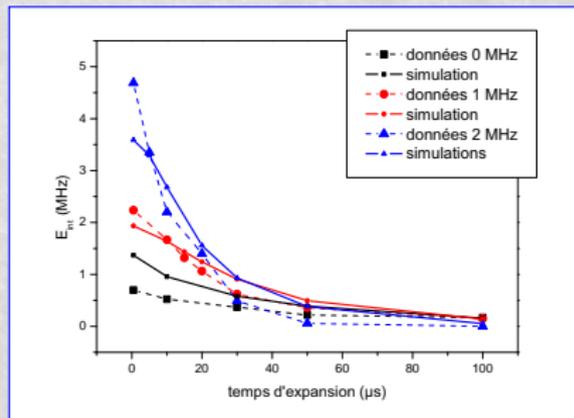
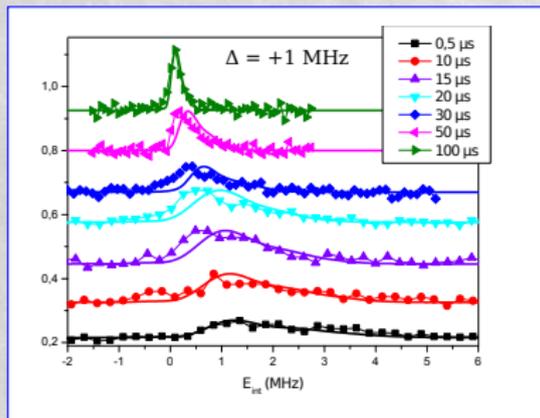


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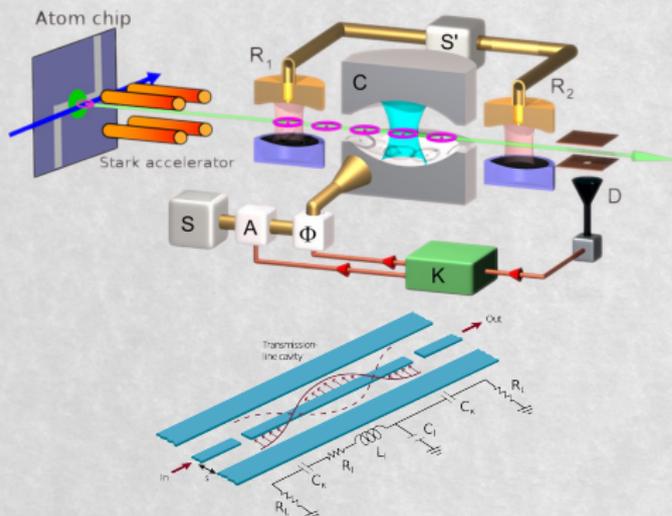
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LONG TERM AIMS

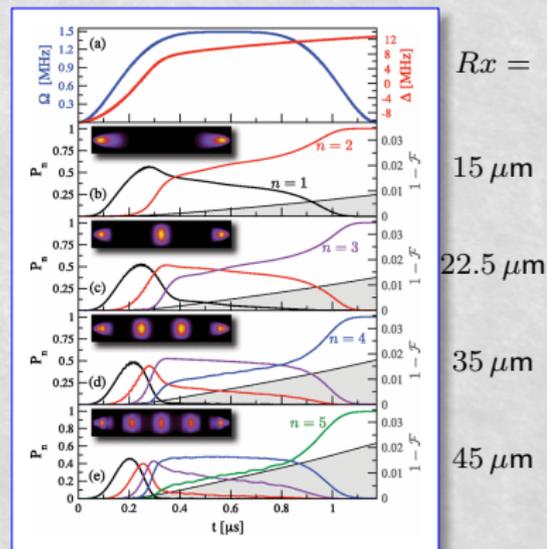


- Adiabatic passage:

Excitation on an elongated BEC, depending on the size of BEC and appropriate detuning
 \rightarrow 1D crystal of Rydberg atoms

T. Pohl *et al.*, PRL **104**, 043002 (2010)

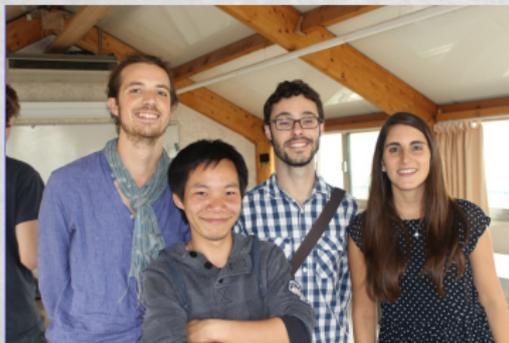
- Couple deterministic single Rydberg atom source and CQED/circuit QED Experiment


 $Rx =$
 $15 \mu\text{m}$
 $22.5 \mu\text{m}$
 $35 \mu\text{m}$
 $45 \mu\text{m}$

TEAM

PhD students

Raul Celistrino Teixeira
Carla Hermann Avigliano
Thanh Long NGUYEN
Tigrane Cantat-Moltrecht



Permanents

Michel Brune
Sébastien Gleyzes
Jean-Michel Raimond
Serge Haroche



Thank you for your attention!

