

# Multimode homodyne detection as a tool for cluster states generation and gaussian quantum computation

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« Gdr Information quantique », 30 Novembre 2012

# Can we use our multimode quantum optical experiment as a quantum computer ?

## → This work :

Theoretical study to guide next future quantum optics experiments  
(N. Treps, C. Fabre, J. Roslund, R. Medeiros de Araujo, P. Jian, Y. Cai)

## → What we have :

Multimode squeezed states of light in a cavity

+

Multipixel homodyne detection system (next future)

## → What we want to do (this talk):

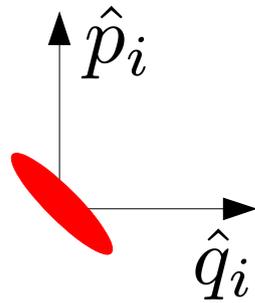
- Demonstrate quadratures measurement with cluster state statistics
- Study the applications to measurement based quantum computation
- ...

# Outlines

- 1 The physical system(s) + what we can do
- 2 Cluster states + an example of implementation
- 3 Measurement based quantum computation + an example of implementation

# Experiment : several co-propagating modes

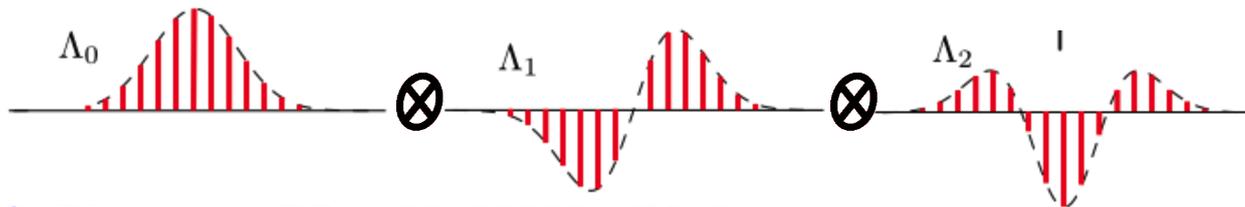
individually squeezed!



Formally described by

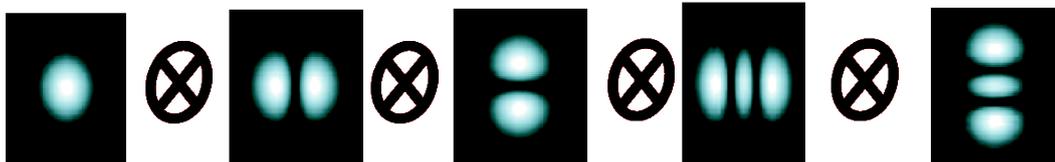
$$\vec{a}_{\text{squ}} = \begin{pmatrix} \hat{a}_{\text{squ}_1} \\ \hat{a}_{\text{squ}_2} \\ \dots \\ \hat{a}_{\text{squ}_N} \end{pmatrix}$$

Either **temporal**...



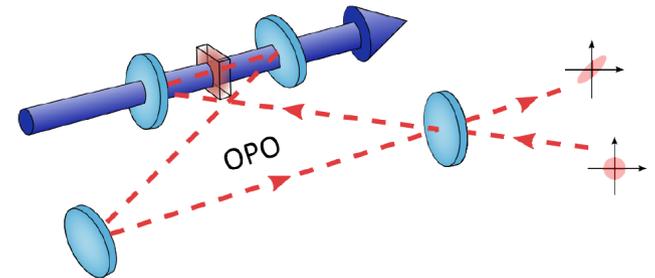
*O. Pinel et al, PRL 108, 083601 (2012) ;  
G. Patera et al, EPJD 56, 123 (2012) PARIS*

...or **spatial**.



*S. Armstrong et al, Nat. Comm. 3, 1026 (2012) CANBERRA*

femto-second laser pump  
+ non-linear crystal  
+ cavity (SPOPO)

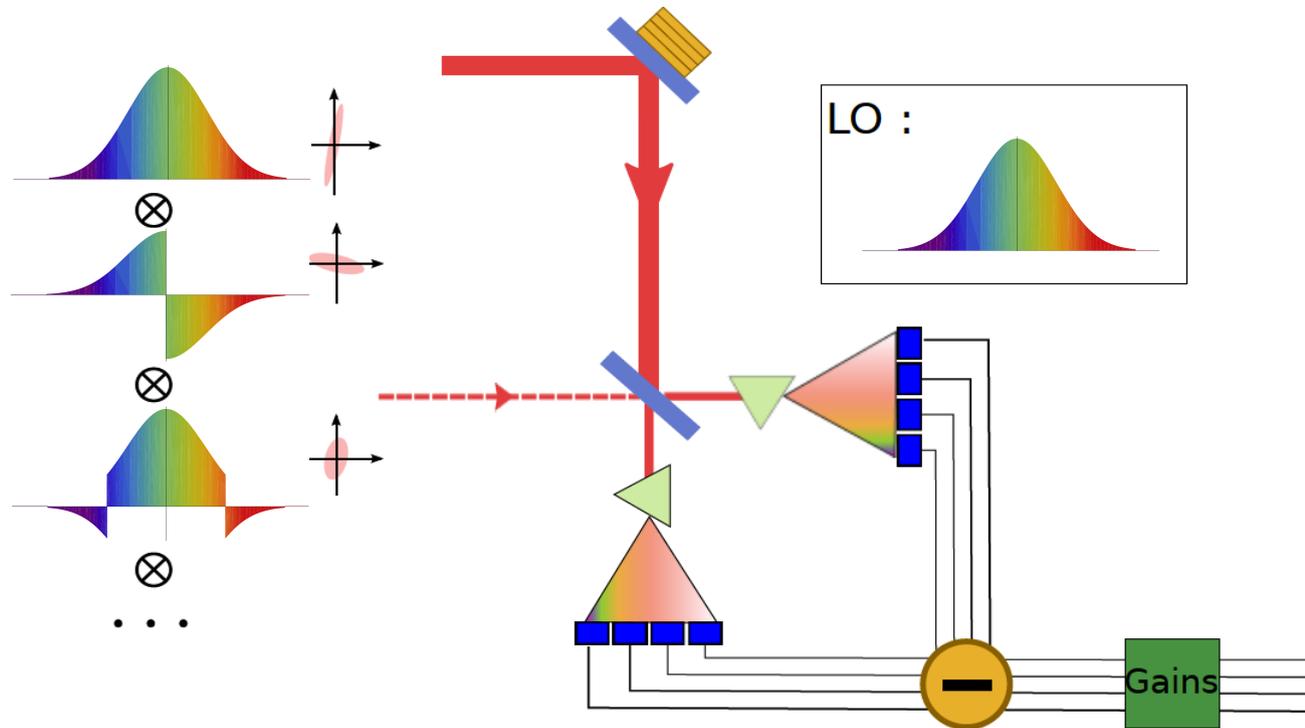


laser pumps  
+ non-linear crystals  
+ several cavities, then  
combined (OPOs)

# MultiPixel Homodyne Detection (MPHD)

Temporal modes

$$\hat{a}_{\text{squ}_i}$$



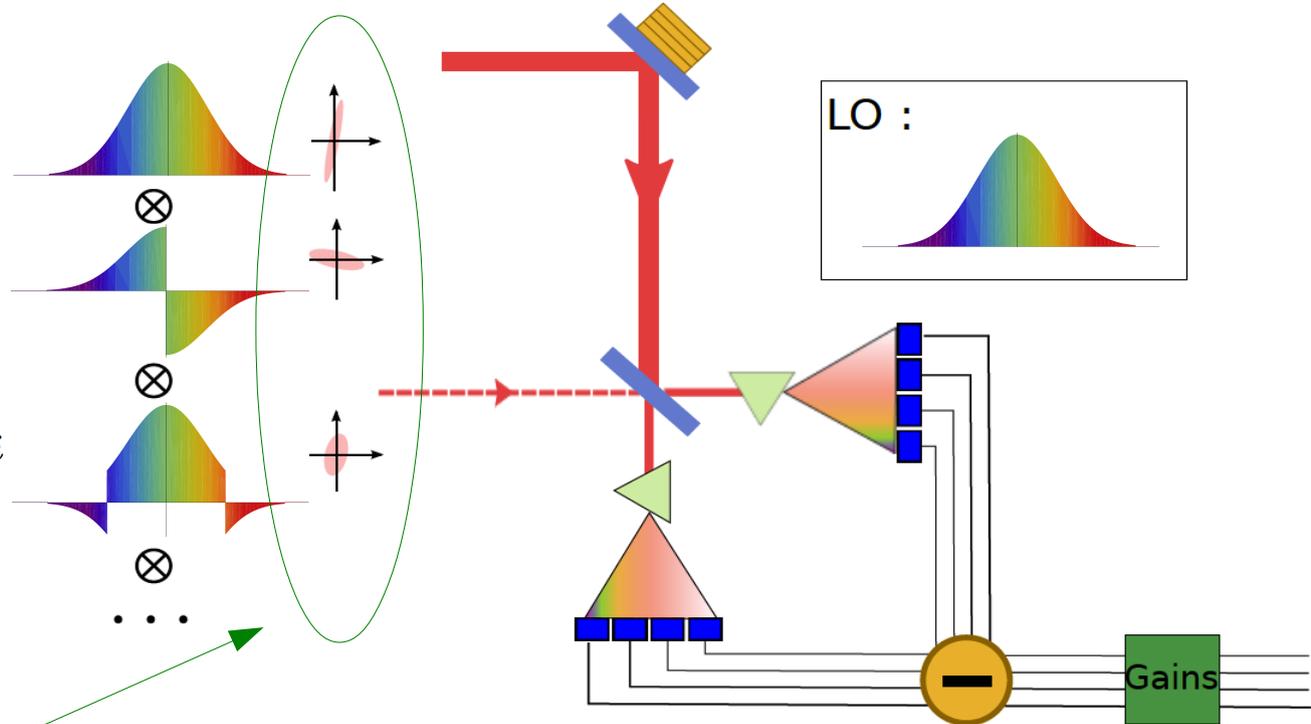
# MultiPixel Homodyne Detection (MPHD)

Temporal modes

$$\hat{a}_{\text{squ}_i}$$

1) input dephasing

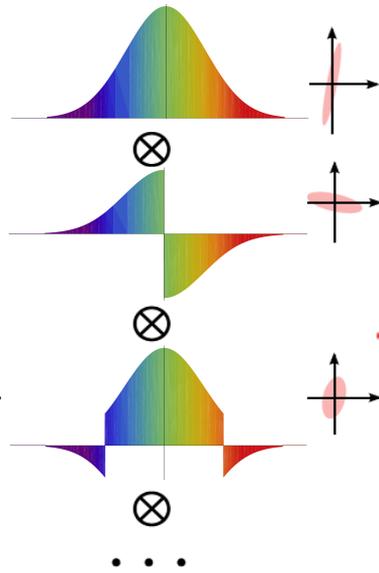
$$\hat{a}_{\text{squ}'_i} = e^{i\varphi_i} \hat{a}_{\text{squ}_i}$$
$$D_r$$



# MultiPixel Homodyne Detection (MPHD)

Temporal modes

$$\hat{a}_{\text{squ}_i}$$

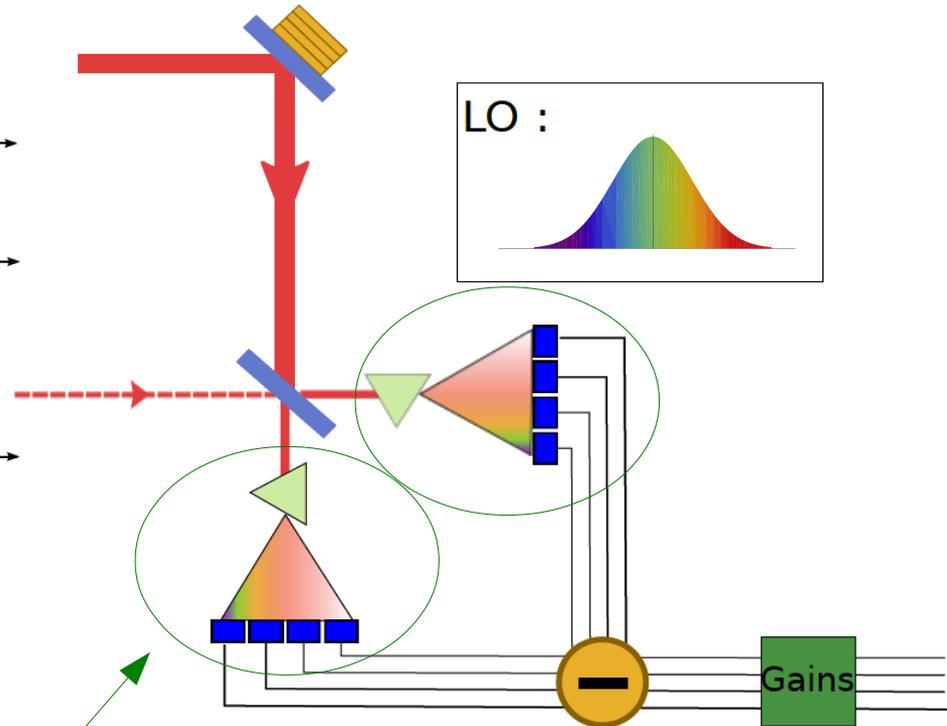
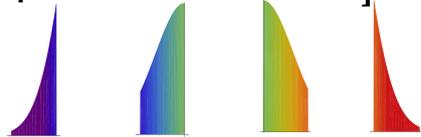


1) input dephasing

$$\hat{a}_{\text{squ}'_i} = e^{i\varphi_i} \hat{a}_{\text{squ}_i} \\ D_r$$

2) detection in the pixel basis

$$\vec{a}_{\text{pix}} = U_T \vec{a}_{\text{squ}'}$$



# Detection matrix (example : 4 modes)

$$\begin{pmatrix} \text{purple peak} \\ \text{blue peak} \\ \text{green peak} \\ \text{red peak} \end{pmatrix} = \kappa \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & -1 & -1 \\ -1 & 1 & 1 & -1 \\ 1 & -1 & 1 & -1 \end{pmatrix}^{-1} \begin{pmatrix} \text{rainbow peak} \\ \text{blue/red split} \\ \text{blue/red inverted split} \\ \text{blue/red split} \end{pmatrix}$$

output of the detectors :  
**pixel modes**

detection matrix  $U_T$   
 ( $\kappa$  normalization constant)

input of the detectors :  
**flip modes**

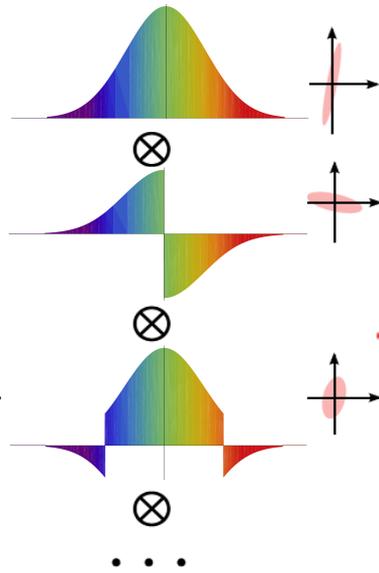
$$I_-^{\varphi_0} \propto \hat{q}_{\text{pix}_i}^{\varphi_0} = (\hat{q}_{\text{pix}_i} \cos \varphi_0 - \hat{p}_{\text{pix}_i} \sin \varphi_0)$$

$\varphi_0$  global phase of the local oscillator

# MultiPixel Homodyne Detection (MPHD)

Temporal modes

$$\hat{a}_{\text{squ}_i}$$

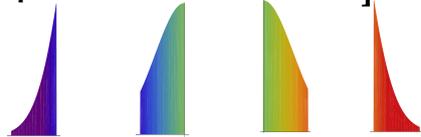


1) input dephasing

$$\hat{a}_{\text{squ}'_i} = e^{i\phi_i} \hat{a}_{\text{squ}_i} D_r$$

2) detection in the pixel basis

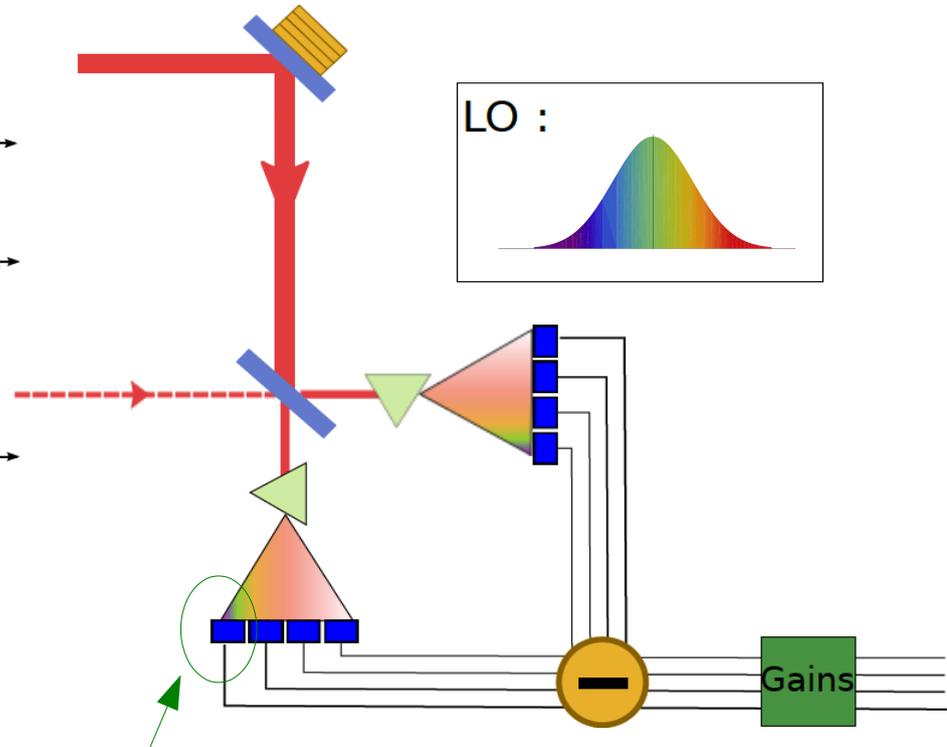
$$\vec{a}_{\text{pix}} = U_T \vec{a}_{\text{squ}'}$$



3) pixel dephasing

$$\hat{a}_{\text{pix}'_i} = e^{i\phi_i} \hat{a}_{\text{pix}_i}$$

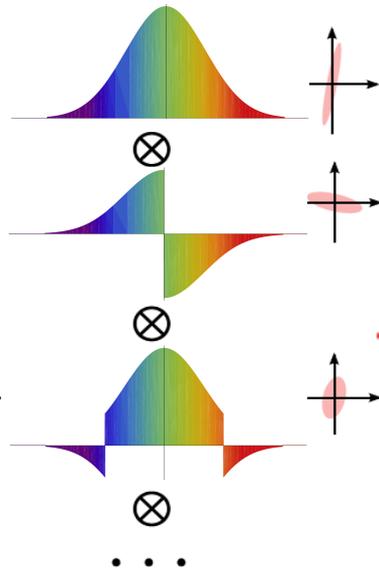
$$\Delta_{\text{LO}}$$



# MultiPixel Homodyne Detection (MPHD)

Temporal modes

$$\hat{a}_{\text{squ}_i}$$

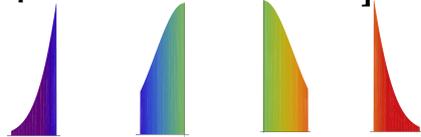


1) input dephasing

$$\hat{a}_{\text{squ}'_i} = e^{i\phi_i} \hat{a}_{\text{squ}_i} D_r$$

2) detection in the pixel basis

$$\vec{a}_{\text{pix}} = U_T \vec{a}_{\text{squ}'}$$



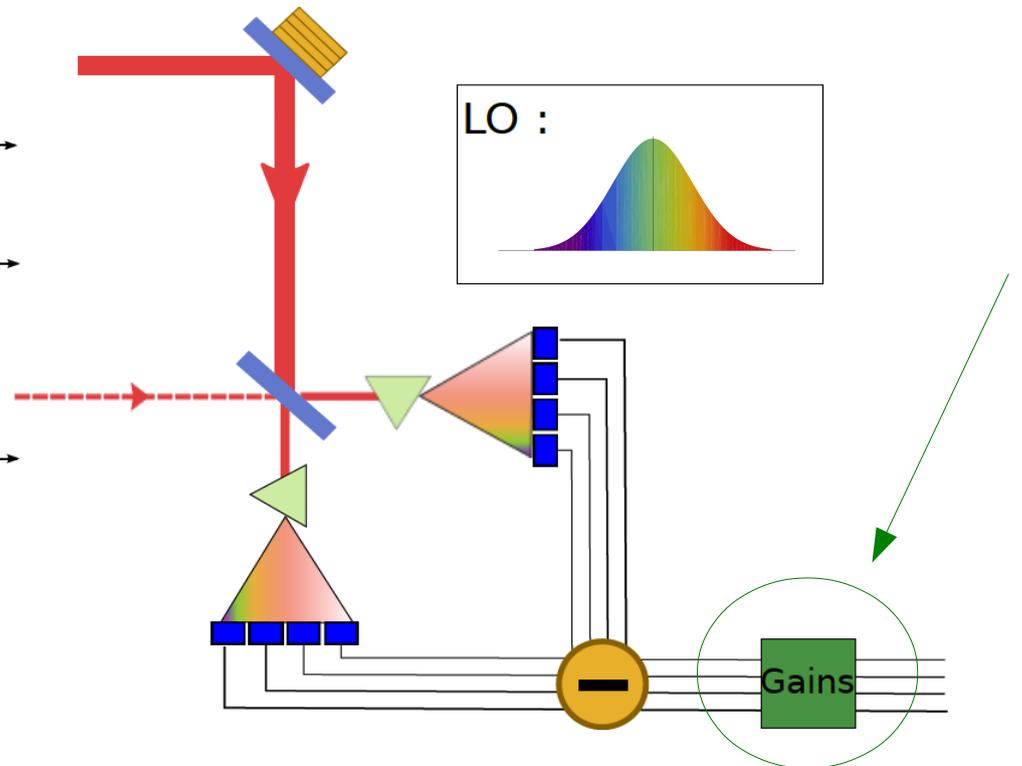
3) pixel dephasing

$$\hat{a}_{\text{pix}'_i} = e^{i\phi_i} \hat{a}_{\text{pix}_i}$$

$$\Delta_{\text{LO}}$$

4) digital recombination

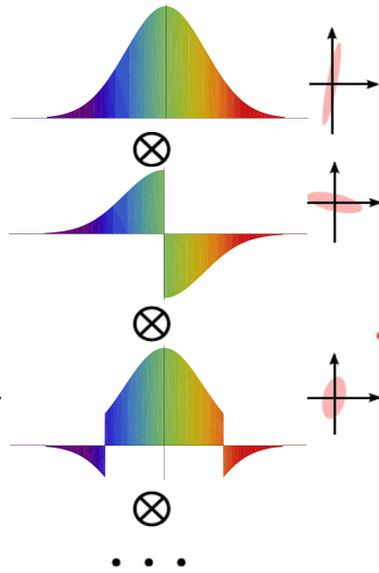
$$\vec{a}_{\text{out}} = O \vec{a}_{\text{pix}'}$$



# MultiPixel Homodyne Detection (MPHD)

Temporal modes

$$\hat{a}_{\text{squ}_i}$$



1) input dephasing

$$\hat{a}_{\text{squ}'_i} = e^{i\varphi_i} \hat{a}_{\text{squ}_i} D_r$$

2) detection in the pixel basis

$$\vec{a}_{\text{pix}} = U_T \vec{a}_{\text{squ}'}$$

3) pixel dephasing

$$\hat{a}_{\text{pix}'_i} = e^{i\phi_i} \hat{a}_{\text{pix}_i} \Delta_{\text{LO}}$$

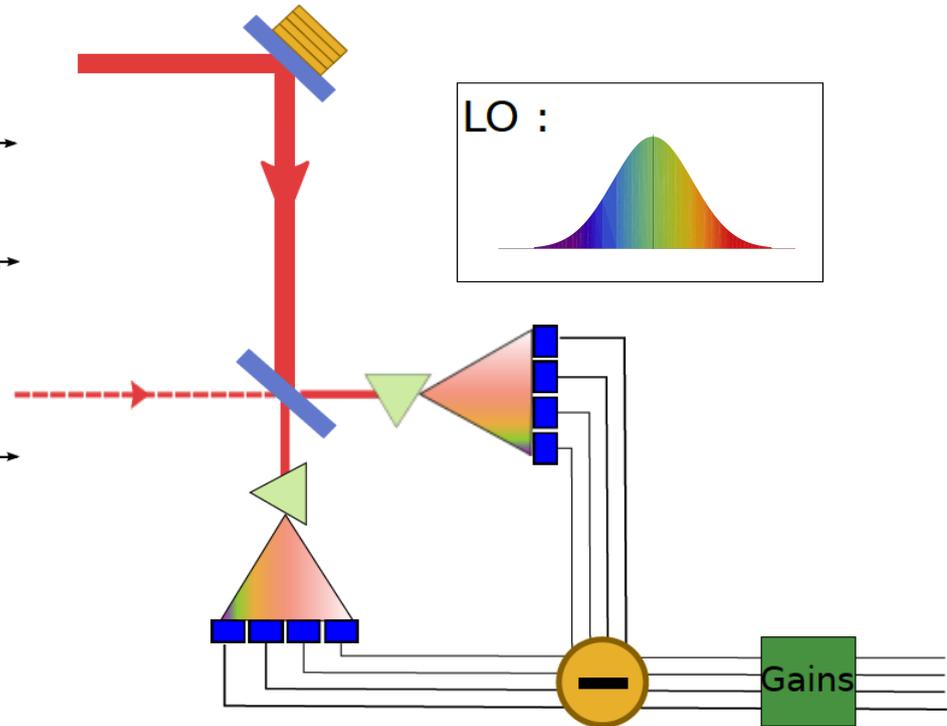
4) digital recombination

$$\vec{a}_{\text{out}} = O \vec{a}_{\text{pix}'}$$

→  $\hat{p}_{\text{out}_i}$

$$\vec{a}_{\text{out}} = O \Delta_{\text{LO}} U_T D_r \vec{a}_{\text{squ}} \equiv U_{\text{MPHD}} \vec{a}_{\text{squ}}$$

UNITARY MATRIX

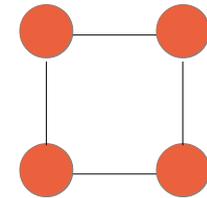


# Why is this interesting? Cluster states :

Operational definition :

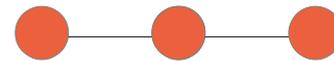
- Take  $N$  (infinitely) squeezed modes

- Apply Cz gates according to graph matrix  $V$



e.g. :

$$V = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$



$$C_Z(|\psi\rangle_i \otimes |\psi\rangle_j) = e^{i\hat{q}_i \otimes \hat{q}_j} (|\psi\rangle_i \otimes |\psi\rangle_j)$$

Symplectic description :

$$\begin{pmatrix} \vec{a}_{\text{clu}} \\ \vec{a}_{\text{clu}}^\dagger \end{pmatrix} = \begin{pmatrix} I & 0 \\ V & I \end{pmatrix} \begin{pmatrix} \vec{a}_{\text{squ}} \\ \vec{a}_{\text{squ}}^\dagger \end{pmatrix}$$

due to Bloch-Messiah decomposition :

$$\begin{pmatrix} \vec{a}_{\text{clu}} \\ \vec{a}_{\text{clu}}^\dagger \end{pmatrix} = \begin{pmatrix} U & 0 \\ 0 & U^* \end{pmatrix} \begin{pmatrix} \vec{a}_{\text{squ}} \\ \vec{a}_{\text{squ}}^\dagger \end{pmatrix}$$

UNITARY MATRIX

*P. Van Loock et al, PRA 76, 032031 (2007) ;*

➔ If  $U_{\text{MPHD}} = U$  we can mimic cluster states statistics with our detection system!

# Which transformations can we implement?

Given a certain transformation  $U$ , can I implement it in the lab, i.e.

$$\exists U_{\text{MPHD}} \text{ such that } U = U_{\text{MPHD}}?$$

YES !  $\longleftrightarrow$   $U'^T U' = D$  with  $U' = U(U_T D_r)^\dagger$   
 $D$  any diagonal matrix with unitary modulus elements

Necessity :

$$U = U_{\text{MPHD}} = O \Delta_{\text{LO}} U_T D_r$$

$$\longrightarrow U' = O \Delta_{\text{LO}} \text{ and } U'^T U' = \Delta_{\text{LO}}^T O^T O \Delta_{\text{LO}} = \Delta_{\text{LO}}^2 = D$$

Sufficiency :

$$U'^T U' = D \longrightarrow \begin{cases} \Delta_{\text{LO}} = D^{\frac{1}{2}} \\ O = U' \Delta_{\text{LO}}^{-1} \end{cases} \text{ used in practice to find the experimental parameter to used!}$$

# An example : four-mode linear cluster state

$$\vec{a}_{\text{clu}} = U_{\text{lin}} \vec{a}_{\text{squ}}$$



$$\text{with } U_{\text{lin}} = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{10}} & \frac{2i}{\sqrt{10}} & 0 \\ \frac{i}{\sqrt{2}} & -\frac{i}{\sqrt{10}} & \frac{2}{\sqrt{10}} & 0 \\ 0 & -\frac{2}{\sqrt{10}} & \frac{i}{\sqrt{10}} & \frac{i}{\sqrt{2}} \\ 0 & -\frac{2i}{\sqrt{10}} & -\frac{1}{\sqrt{10}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

We can find a solution for the experimental parameters

$$U_{\text{lin}} = U_{\text{MPHD}} = O \Delta_{\text{LO}} U_T D_r$$

$$\gamma \equiv \arctan(1/2)$$

$$\Delta_{\text{LO}} = \begin{pmatrix} e^{i\frac{\gamma+\pi}{2}} & 0 & 0 & 0 \\ 0 & e^{-i\frac{\gamma}{2}} & 0 & 0 \\ 0 & 0 & e^{i\frac{\gamma}{2}} & 0 \\ 0 & 0 & 0 & e^{-i\frac{\gamma+\pi}{2}} \end{pmatrix} \quad O = \begin{pmatrix} -\frac{1}{2}\sqrt{1-\frac{2}{\sqrt{5}}} & \frac{1}{2}\sqrt{1+\frac{2}{\sqrt{5}}} & \sqrt{\frac{1}{4}+\frac{1}{2\sqrt{5}}} & -\frac{1}{2}\sqrt{1-\frac{2}{\sqrt{5}}} \\ \sqrt{\frac{1}{4}+\frac{1}{2\sqrt{5}}} & -\frac{1}{2}\sqrt{1-\frac{2}{\sqrt{5}}} & \frac{1}{2}\sqrt{1-\frac{2}{\sqrt{5}}} & -\frac{1}{2}\sqrt{1+\frac{2}{\sqrt{5}}} \\ \sqrt{\frac{1}{4}+\frac{1}{2\sqrt{5}}} & \frac{1}{2}\sqrt{1-\frac{2}{\sqrt{5}}} & \frac{1}{2}\sqrt{1-\frac{2}{\sqrt{5}}} & \frac{1}{2}\sqrt{1+\frac{2}{\sqrt{5}}} \\ -\frac{1}{2}\sqrt{1-\frac{2}{\sqrt{5}}} & -\frac{1}{2}\sqrt{1+\frac{2}{\sqrt{5}}} & \sqrt{\frac{1}{4}+\frac{1}{2\sqrt{5}}} & \frac{1}{2}\sqrt{1-\frac{2}{\sqrt{5}}} \end{pmatrix}$$

$$U_T \text{ fixed detection matrix ; } \quad D_r = \text{diag}(1, i, i, 1)$$

# So in practice we have to set...

Temporal modes

$$a_{\text{squ}_i}$$

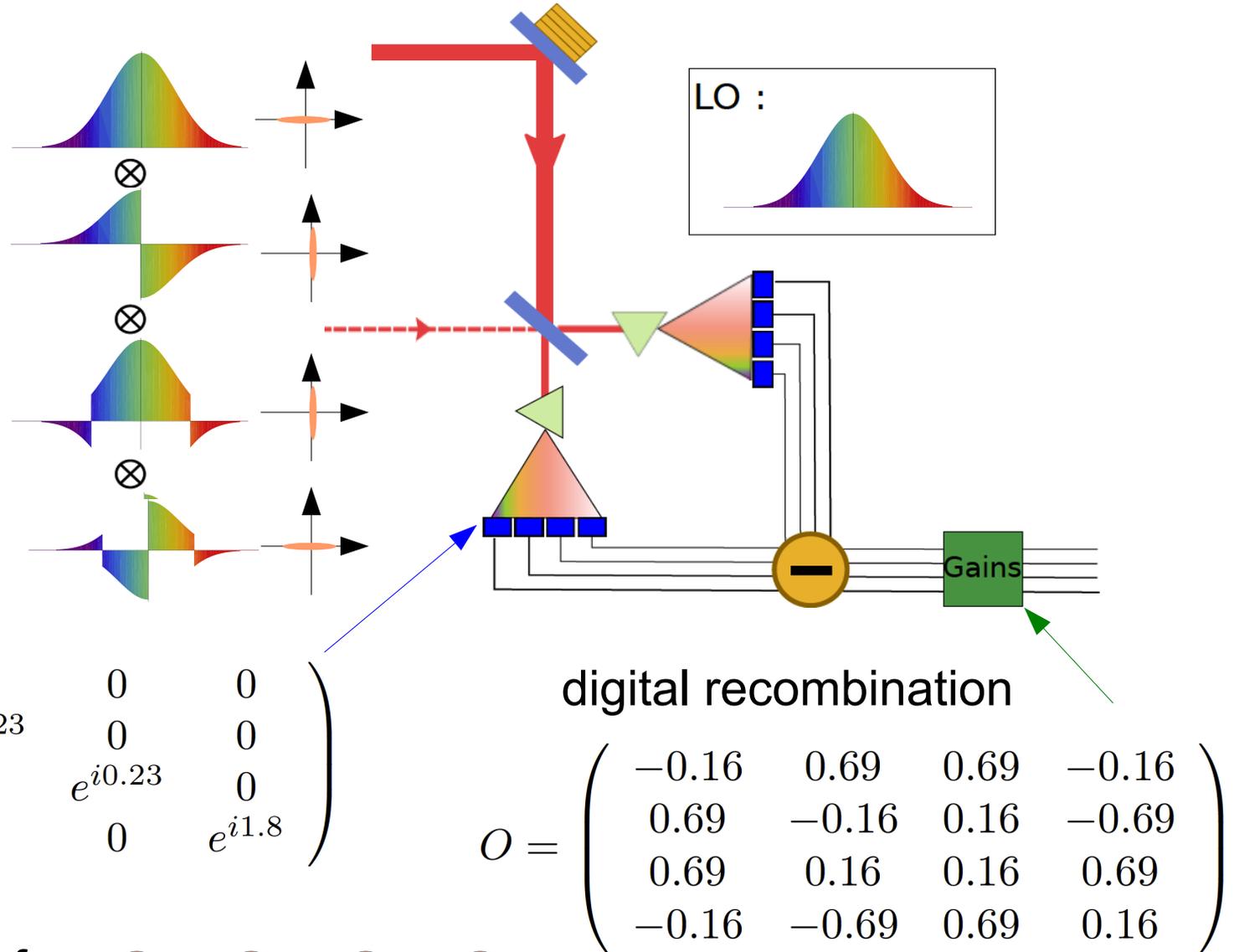
input dephasing

$$D_r = \text{diag}(1, i, i, 1)$$

pixel dephasing

$$\Delta_{\text{LO}} = \begin{pmatrix} e^{i1.8} & 0 & 0 & 0 \\ 0 & e^{-i0.23} & 0 & 0 \\ 0 & 0 & e^{i0.23} & 0 \\ 0 & 0 & 0 & e^{i1.8} \end{pmatrix}$$

→  $\hat{p}_{\text{out}_i}$  as for 



# CV quantum computation

## Definition :

« A system is a universal quantum computer if it can simulate the action of a Hamiltonian consisting of a general polynomial of  $\hat{q}$  and  $\hat{p}$  to any fixed accuracy »

$$|\psi\rangle_{\text{out}} = e^{iH(\hat{q},\hat{p})t} |\psi\rangle_{\text{in}} \quad \text{M. Gu et al, PRA 79, 062318 (2009)}$$

How is this reached ?

→ **Universal gate set** :  $\left\{ D_{f(\hat{q})}, e^{i\hat{q}s}, e^{i\hat{q}^2 s}, e^{i\frac{\pi}{4}(\hat{q}^2 + \hat{p}^2)}, e^{i\hat{q}_1 \otimes \hat{q}_2}, e^{i\hat{q}^3 s} \right\}$

↑
↑
↑

Single-mode gaussian
Two-mode
non-gaussian

→ **Short time evolution (Trotter expansion)**

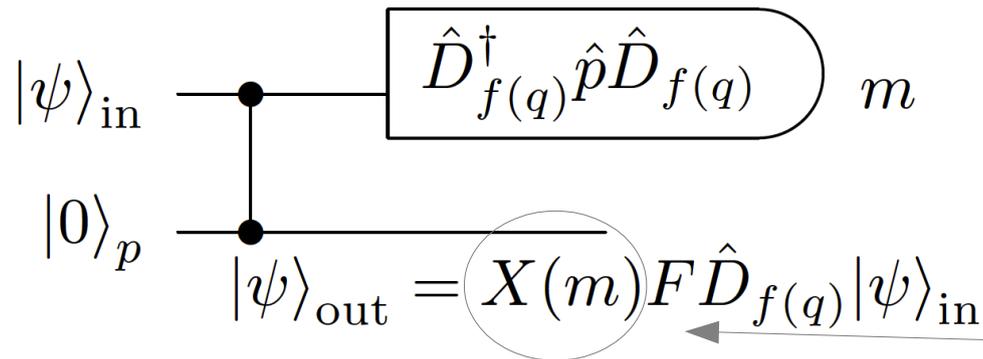
$$e^{iA\delta t} e^{iB\delta t} e^{-iA\delta t} e^{-iB\delta t} = e^{-[A,B]\delta t^2 + O(\delta t^3)}$$

$$\{H_i\} \Rightarrow \pm i[H_i, H_j], \pm [H_i, [H_j, H_k]] \dots$$

S. Lloyd and S. Braunstein, PRL 82, 1784 (1999)

# Measurement based quantum computation

I want to implement the operation  $D_{f(\hat{q})}$ :

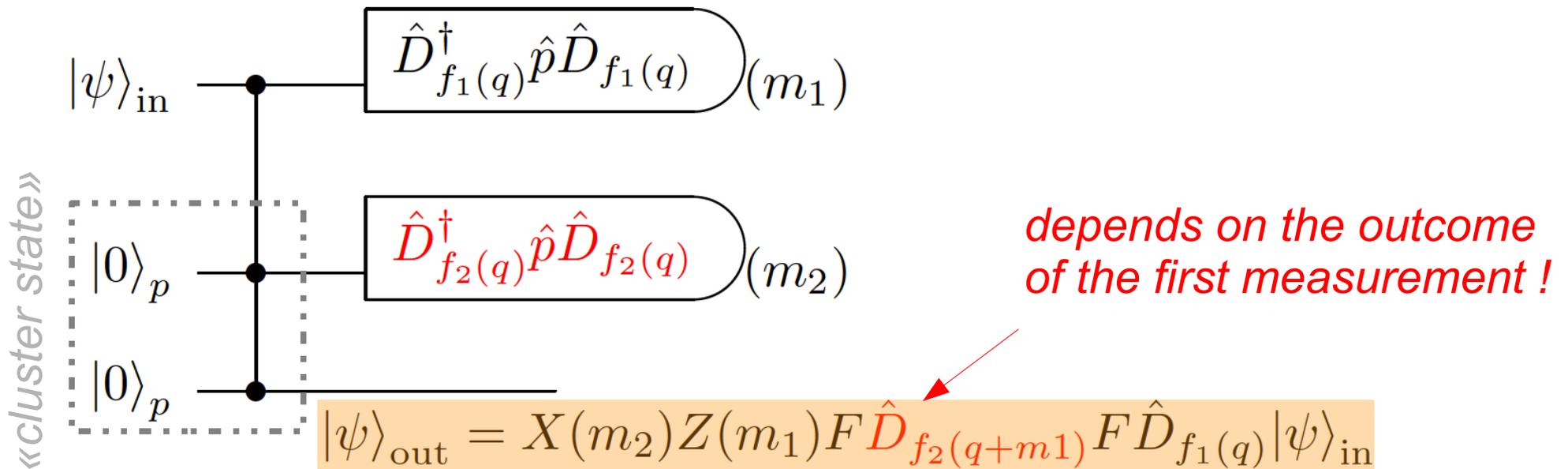


with  $\begin{array}{c} \bullet \\ | \\ \bullet \end{array} \equiv e^{i\hat{q}_1 \otimes \hat{q}_2}$   
(Cz gate)

and  $X(s) = e^{-is\hat{p}}$

« by-product operator »

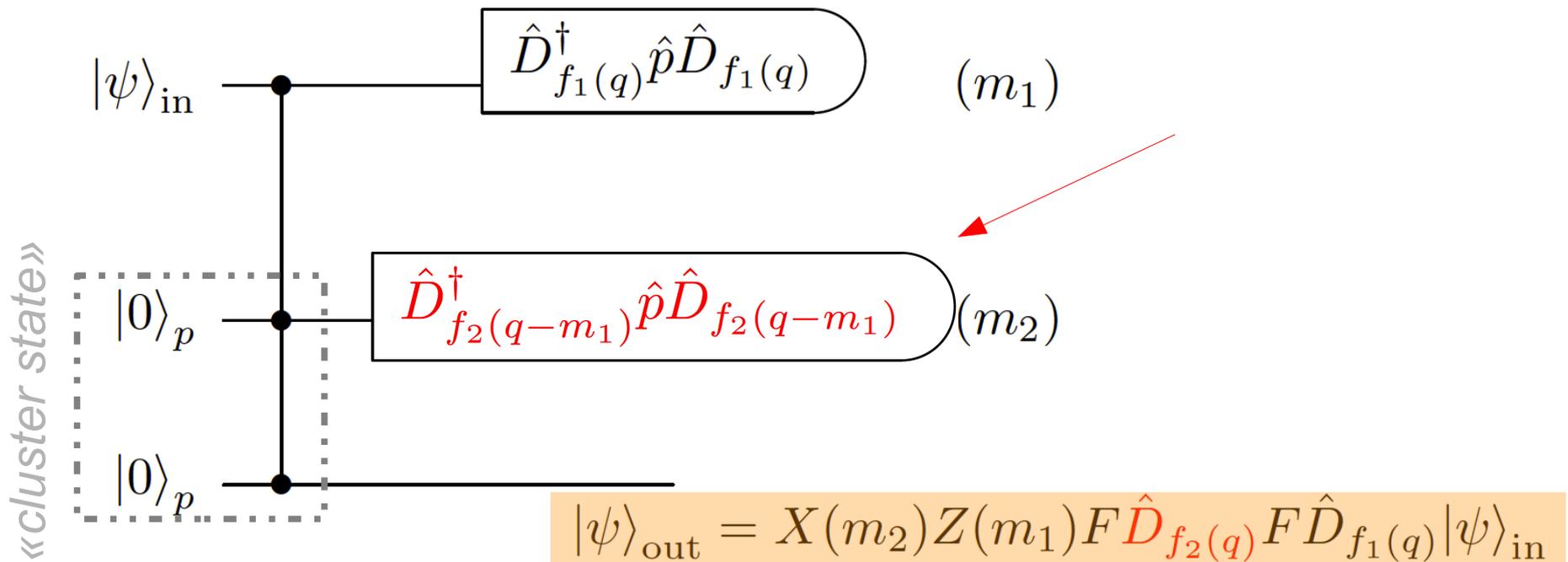
For more operations, repeat the previous circuit :



*depends on the outcome of the first measurement !*

# Measurement based quantum computation

To obtain *deterministic operations* one has to *adapt the measurement basis depending on the outcomes of previous measurements*



# Implementation of Gaussian unitaries

→ by homodyne detection ; e.g :  $D_{f(\hat{q})} = e^{is\hat{q}^2}$

$$e^{-is\hat{q}^2} \hat{p} e^{is\hat{q}^2} = \hat{p} + s\hat{q} = g(\hat{q} \sin \theta + \hat{p} \cos \theta)$$

with  $g = \sqrt{1 + s^2}$  and  $\theta = \arctan s$

If I have to adapt my measurement basis :

$$e^{-is(\hat{q}+m)^2} \hat{p} e^{is(\hat{q}+m)^2} = \hat{p} + s\hat{q} + ms$$

Can be done while post-processing!



**ALL GAUSSIAN MEASUREMENTS CAN BE PERFORMED SIMULTANEOUSLY («PARALLELISM»)**

*M. Gu et al, PRA 79, 062318 (2009)*

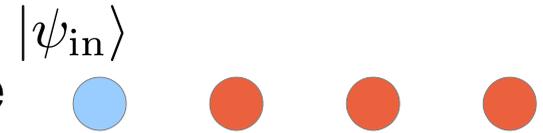
This is not the case for : **cubic phase gate**  $D_{f(\hat{q})} = e^{i\gamma \frac{\hat{q}^3}{3}}$

*D. Gottesman et al, PRA 64, 012310 (2001)*

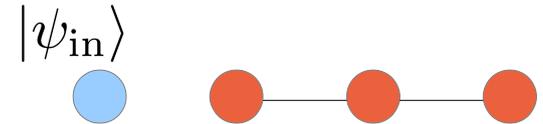
# A simple example : Fourier transform

$$\begin{pmatrix} \hat{q}_{\text{out}} \\ \hat{p}_{\text{out}} \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \hat{q}_{\text{in}} \\ \hat{p}_{\text{in}} \end{pmatrix} = \begin{pmatrix} -\hat{p}_{\text{in}} \\ \hat{q}_{\text{in}} \end{pmatrix} \quad \text{How to realize that in MBQC?}$$

1) Start with 3 squeezed modes and an input mode

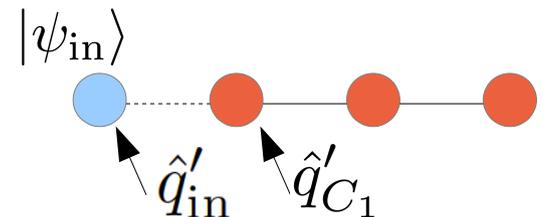


2) Create a 3-mode cluster state  $\vec{a}_C = U_{\text{CL}}^{1,2,3} \vec{a}_{\text{squ}}$



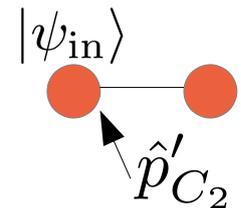
3) Mix the input state to the 1<sup>st</sup> mode of the cluster with a beam splitter

$$\begin{pmatrix} \hat{a}'_{\text{in}} \\ \hat{a}'_{C_1} \end{pmatrix} = U_{\text{BS}}^{\text{in}, C_1} \begin{pmatrix} \hat{a}_{\text{in}} \\ \hat{a}_1 \end{pmatrix}$$



4) Measure  $\hat{q}'_{\text{in}}$  and  $\hat{q}'_{C_1}$  to teleport the input state

5) Measure  $\hat{p}'_{C_2}$  on the 2<sup>nd</sup> mode of the cluster



6) Read out in arbitrary direction

$$|\psi_{\text{out}}\rangle = F|\psi_{\text{in}}\rangle$$

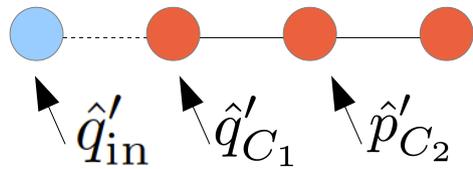
A single red circle is positioned below the equation.

# Fourier transform of an input state

*Can we implement it with our experiment ?*

Starting from 3 squeezed modes + an input mode to be processed :

« Theory » :



$$U_{\text{FT}} = D_{\text{meas}}(U_{\text{BS}}^{\text{in},1} \otimes \mathcal{I}^{2,3})(\mathcal{I}^{\text{in}} \otimes U_{\text{CL}}^{1,2,3})$$

$$\theta_3 = 0 \quad = \begin{pmatrix} \frac{i}{\sqrt{2}} & 0 & \frac{1}{\sqrt{3}} & \frac{i}{\sqrt{6}} \\ -\frac{1}{\sqrt{2}} & 0 & -\frac{i}{\sqrt{3}} & \frac{1}{\sqrt{6}} \\ 0 & -\frac{i}{\sqrt{2}} & -\frac{i}{\sqrt{6}} & -\frac{1}{\sqrt{3}} \\ 0 & -\frac{1}{\sqrt{2}} & \frac{1}{\sqrt{6}} & -\frac{i}{\sqrt{3}} \end{pmatrix}$$

« Experiment » :  $U_{\text{MPHD}} = O\Delta_{\text{LO}}U_T D_r$

**CAN WE SET  $U_{\text{MPHD}} = U_{\text{FT}}$ ?**



LET'S CHECK WITH  
OUR METHOD !

# Solution for the Fourier Transform

Compute  $U'^T U' = D$ , then

$$\Delta_{\text{LO}} = D^{\frac{1}{2}} \quad O = U' \Delta_{\text{LO}}^{-1}$$

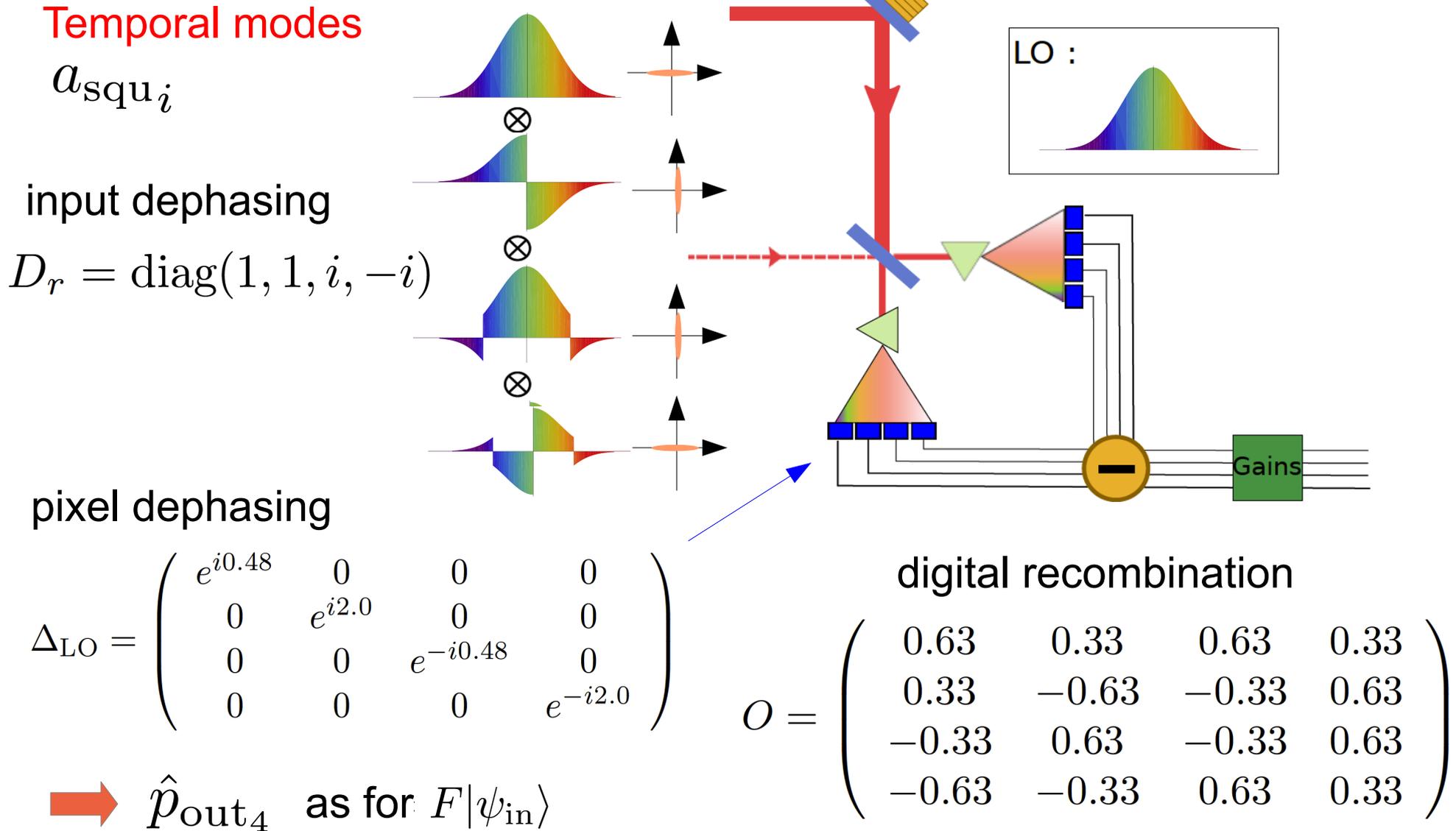
$$O = \begin{pmatrix} \frac{1}{2} \sqrt{1 + \sqrt{\frac{2}{3}}} & \frac{1}{2\sqrt{3+\sqrt{6}}} & -\frac{1}{2\sqrt{3+\sqrt{6}}} & -\frac{1}{2} \sqrt{1 + \sqrt{\frac{2}{3}}} \\ \frac{1}{2\sqrt{3+\sqrt{6}}} & -\frac{1}{2} \sqrt{1 + \sqrt{\frac{2}{3}}} & -\frac{1}{2} \sqrt{1 + \sqrt{\frac{2}{3}}} & \frac{1}{2\sqrt{3+\sqrt{6}}} \\ -\frac{1}{2} \sqrt{1 + \sqrt{\frac{2}{3}}} & -\frac{1}{2\sqrt{3+\sqrt{6}}} & -\frac{1}{2\sqrt{3+\sqrt{6}}} & -\frac{1}{2} \sqrt{1 + \sqrt{\frac{2}{3}}} \\ \frac{1}{2\sqrt{3+\sqrt{6}}} & -\frac{1}{2} \sqrt{1 + \sqrt{\frac{2}{3}}} & \frac{1}{2} \sqrt{1 + \sqrt{\frac{2}{3}}} & -\frac{1}{2\sqrt{3+\sqrt{6}}} \end{pmatrix}$$

$$\Delta_{\text{LO}} = \begin{pmatrix} e^{i\frac{(\beta+\pi)}{2}} & 0 & 0 & 0 \\ 0 & e^{i\frac{\beta}{2}} & 0 & 0 \\ 0 & 0 & e^{-i\frac{\beta}{2}} & 0 \\ 0 & 0 & 0 & e^{-i\frac{(\beta+\pi)}{2}} \end{pmatrix} \quad \beta = \arctan 1/\sqrt{2}$$

$$D_r = \text{diag}(1, 1, i, -i)$$

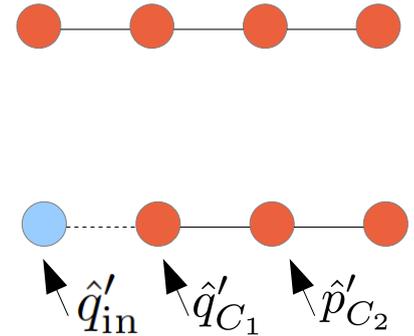
To tune the dephasings: act on pump shape or on the length of the crystal !

# So in practice we have to set...



# Conclusions

- Possibility of measuring cluster state statistics in a compact way
- Possibility of performing gaussian measurement based quantum computation in a compact way



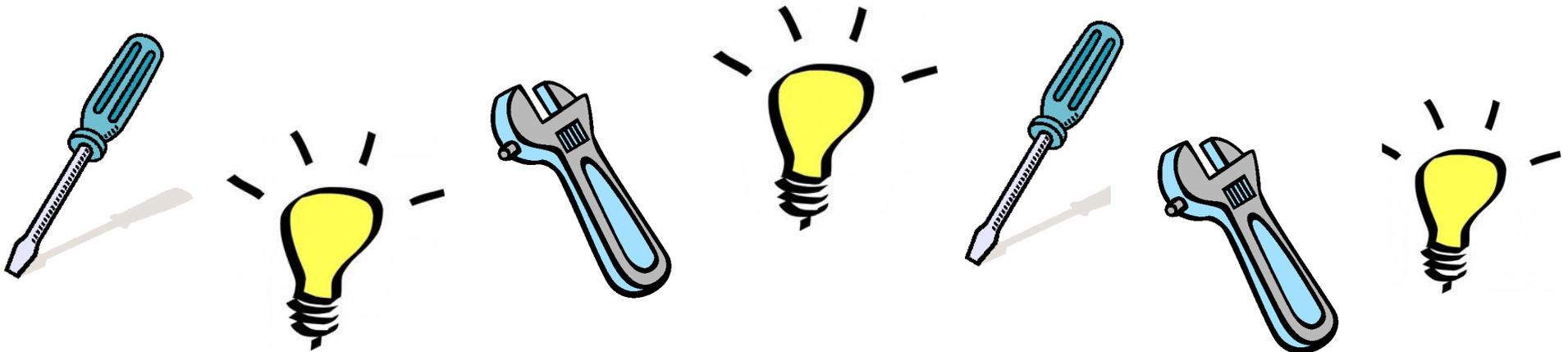
# Perspectives

- Inclusion of finite-squeezing effects (in progress)
- Demonstrate more operations (e.g.: two-mode operations)
- Inclusion of a non-gaussian measurement to go beyond gaussian operations



And then...

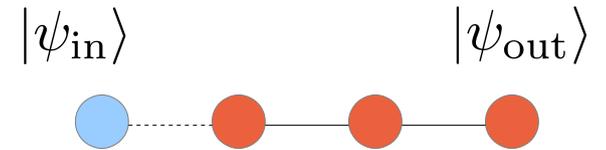
...do the experiment !!!



*THANK YOU FOR YOUR ATTENTION !*

# A simple example : Fourier transform

$$\begin{pmatrix} \hat{q}_{\text{out}} \\ \hat{p}_{\text{out}} \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \hat{q}_{\text{in}} \\ \hat{p}_{\text{in}} \end{pmatrix} = \begin{pmatrix} -\hat{p}_{\text{in}} \\ \hat{q}_{\text{in}} \end{pmatrix}$$



How to realize that in MBQC? Theory :

- Create a 3-mode cluster state

$$\vec{a}_C = U_{\text{CL}}^{1,2,3} \vec{a}_{\text{squ}}$$

$$U_{\text{CL}}^{1,2,3} = \begin{pmatrix} 0 & -\sqrt{\frac{2}{3}} & -\frac{i}{\sqrt{3}} \\ -\frac{i}{\sqrt{2}} & -\frac{i}{\sqrt{6}} & -\frac{1}{\sqrt{3}} \\ -\frac{1}{\sqrt{2}} & \frac{1}{\sqrt{6}} & -\frac{i}{\sqrt{3}} \end{pmatrix}$$

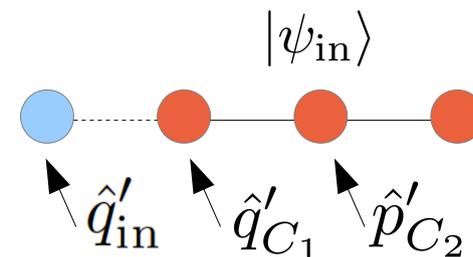
- Attach the input state to the cluster via teleportation (BS + meas.  $\hat{q}'_{\text{in}}, \hat{q}'_1$ )

$$\begin{pmatrix} \hat{a}'_{\text{in}} \\ \hat{a}'_{C_1} \end{pmatrix} = U_{\text{BS}}^{\text{in}, C_1} \begin{pmatrix} \hat{a}_{\text{in}} \\ \hat{a}_1 \end{pmatrix}$$

$$U_{\text{BS}}^{\text{in}, C_1} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & i \\ i & 1 \end{pmatrix}$$

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- Measure  $\hat{p}'_{C2}$  on the node 2 of the cluster; node 3 contains  $|\psi_{\text{out}}\rangle = F|\psi_{\text{in}}\rangle$

Equivalent to applying rotation  $\vec{a}'' = D_{\text{meas}} \vec{a}'$   $D_{\text{meas}} = \text{diag}(e^{i\theta_{\text{in}}}, e^{i\theta_1}, e^{i\theta_2}, e^{i\theta_3})$   
 and measuring  $\hat{p}''_i = (\hat{q}'_i \sin \theta_i + \hat{p}'_i \cos \theta_i)$   $\{\theta_i\} = (\pi/2, \pi/2, 0, \theta_3)$

# Finite squeezing effects

Cluster state generation with a QND interaction

$$S_{K_{\text{EXP}},V} = \begin{pmatrix} I & 0 \\ V & I \end{pmatrix} \begin{pmatrix} K_{\text{EXP}}^{-\frac{1}{2}} & 0 \\ 0 & K_{\text{EXP}}^{\frac{1}{2}} \end{pmatrix}$$

Cluster state generation with linear optics only (+ squeezing off-line)

$$S_{U',V'} = \begin{pmatrix} X_V & -Y_V \\ Y_V & X_V \end{pmatrix} \begin{pmatrix} K_{\text{EXP}}^{-\frac{1}{2}} & 0 \\ 0 & K_{\text{EXP}}^{\frac{1}{2}} \end{pmatrix}$$

# Finite squeezing effects

$$\begin{aligned} \begin{pmatrix} \vec{x}' \\ \vec{p}' \end{pmatrix} &= \begin{pmatrix} X_V & -Y_V \\ Y_V & X_V \end{pmatrix} \begin{pmatrix} K_{\text{EXP}}^{-\frac{1}{2}} & 0 \\ 0 & K_{\text{EXP}}^{\frac{1}{2}} \end{pmatrix} \begin{pmatrix} \vec{x}^{(0)} \\ \vec{p}^{(0)} \end{pmatrix} \\ &= \begin{pmatrix} X_V K_{\text{EXP}}^{-\frac{1}{2}} \vec{x}^{(0)} - Y_V K_{\text{EXP}}^{\frac{1}{2}} \vec{p}^{(0)} \\ Y_V K_{\text{EXP}}^{-\frac{1}{2}} \vec{x}^{(0)} + X_V K_{\text{EXP}}^{\frac{1}{2}} \vec{p}^{(0)} \end{pmatrix}. \end{aligned}$$

$$\vec{p}^C - V \vec{x}^C = 0$$

$$(Y_V - V X_V) K_{\text{EXP}}^{-\frac{1}{2}} \vec{x}^{(0)} = 0 \quad \Rightarrow \quad (Y_V - V X_V) = 0;$$

$$(X_V + V Y_V) K_{\text{EXP}}^{\frac{1}{2}} \vec{p}^{(0)} \rightarrow 0 \quad \text{excess noise.}$$

# Finite squeezing effects

Excess noise quadratures

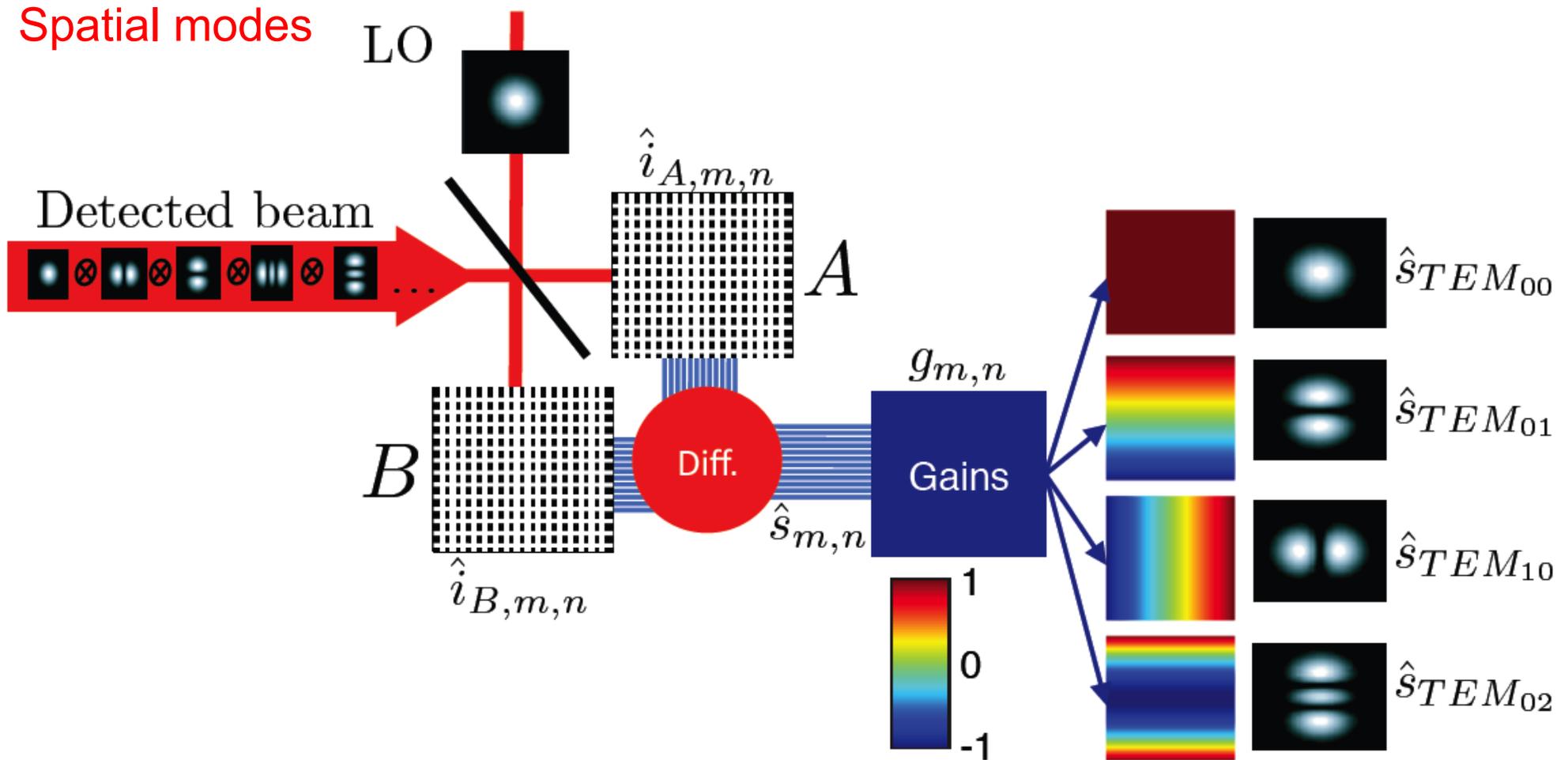
$$\hat{\delta}_1 = \hat{p}'_1 - \sum_l V_{1l} \hat{x}'_l = \hat{p}'_1 - \hat{x}'_2 = \sqrt{2} e^{-r_1} \hat{p}_1^{(0)}$$

$$\hat{\delta}_2 = \hat{p}'_2 - \sum_l V_{2l} \hat{x}'_l = \hat{p}'_2 - \hat{x}'_1 - \hat{x}'_3 = \sqrt{\frac{5}{2}} e^{-r_3} \hat{p}_3^{(0)} + \frac{e^{-r_4}}{\sqrt{2}} \hat{p}_4^{(0)}$$

$$\hat{\delta}_3 = \hat{p}'_3 - \sum_l V_{3l} \hat{x}'_l = \hat{p}'_3 - \hat{x}'_2 - \hat{x}'_4 = \frac{e^{-r_1}}{\sqrt{2}} \hat{p}_1^{(0)} - \sqrt{\frac{5}{2}} e^{-r_2} \hat{p}_2^{(0)}$$

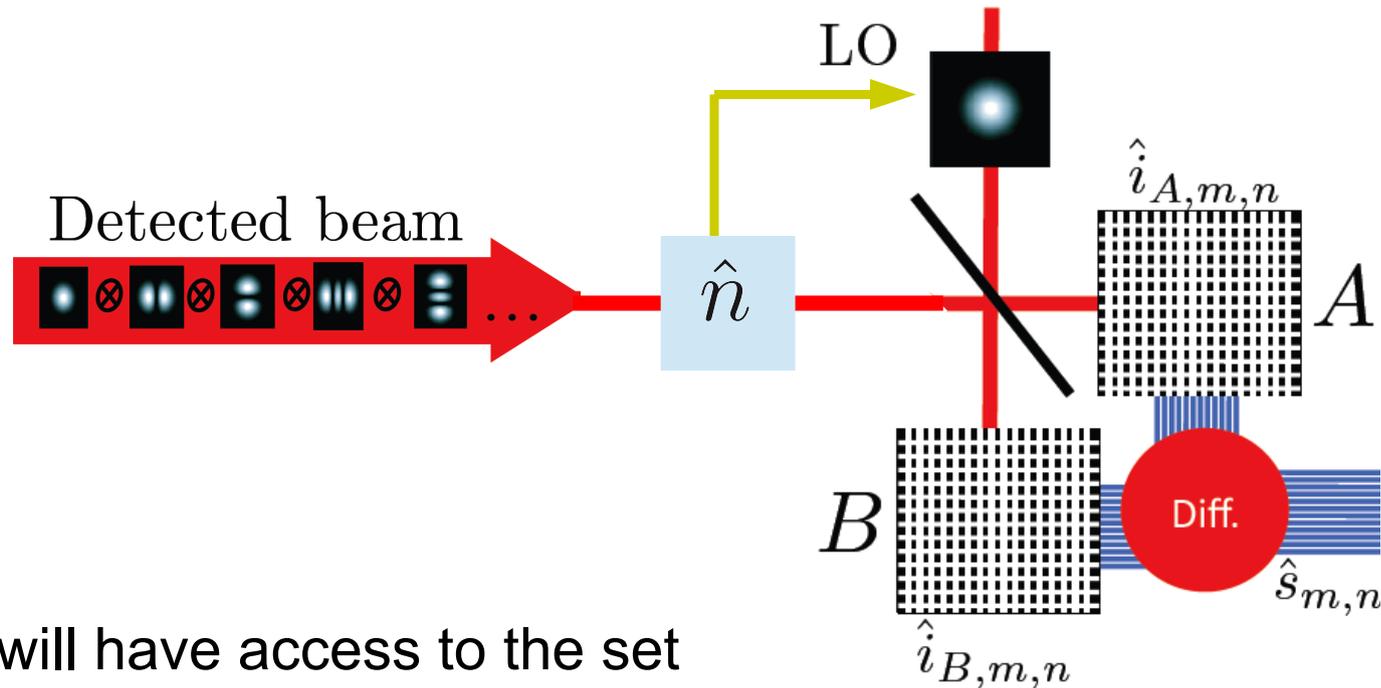
$$\hat{\delta}_4 = \hat{p}'_4 - \sum_l V_{4l} \hat{x}'_l = \hat{p}'_4 - \hat{x}'_3 = \sqrt{2} e^{-r_4} \hat{p}_4^{(0)}$$

# MultiPixel Homodyne Detection (MPHD)



J-F. Morizur, PhD thesis (2011) ; S. Armstrong et al, Nat. Comm.... (2012)

At best, we will be able to...



Hence we will have access to the set

$$e^{i\gamma \frac{\hat{q}^3}{3}} \cdot \{\text{set of accessible gaussian operations}\}$$

- *Which gaussian operations are we able to implement ?*
- *Which total H are we able to implement ?*  
*Do they correspond to any interesting physical phenomenon ?*

# Implementation of non-Gaussian operations

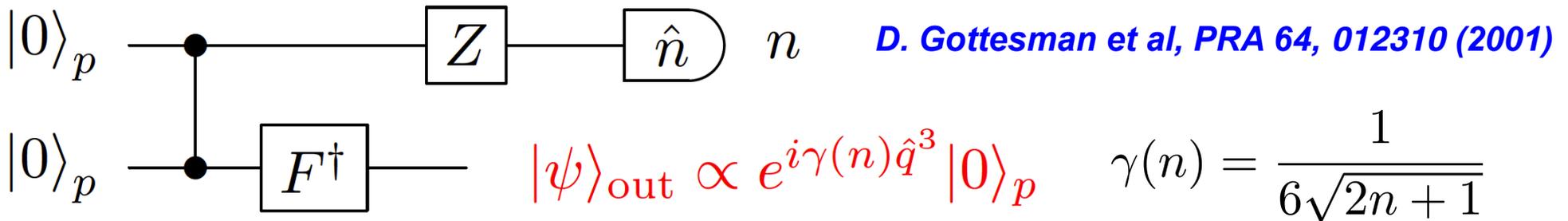
Simplest example : **cubic phase gate**  $D_{f(\hat{q})} = e^{i\gamma \frac{\hat{q}^3}{3}}$

$$\hat{p}_{\gamma(\hat{q}+m)^3/3} = \hat{p} + \gamma\hat{q}^2 + \underbrace{2ms\hat{q}}_{\text{requires actual adaptation...}} + m^2\gamma$$

...which we cannot do!  $\rightarrow$  *at best, we will be able to implement it ONCE !*

Acting the gate on a special initial state :

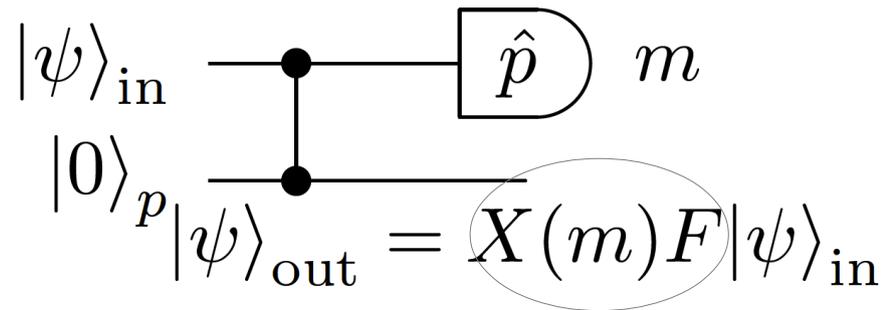
$$Z(s) = e^{is\hat{q}}$$



$\longrightarrow$  **Non-deterministic non-gaussian state**

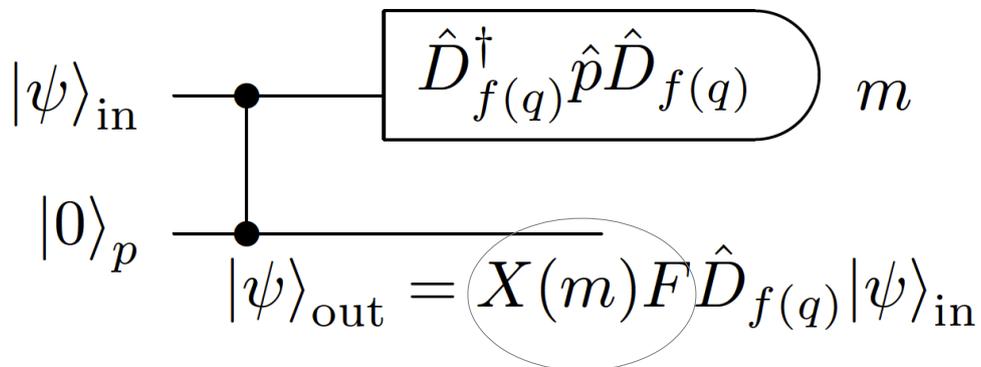
Even for a single first implementation, *adaptation is required for determinism !*

# Measurement based quantum computation



with  $\begin{array}{c} \bullet \\ | \\ \bullet \end{array} \equiv e^{i\hat{q}_1 \otimes \hat{q}_2}$   
 (Cz gate)  
 and  $X(s) = e^{-is\hat{p}}$

« by-product » operators



# How to compute the quadrature to measure?

Several steps of the computation:  $|\psi\rangle_{\text{out}} = \dots F D_{f_2(q)} F D_{f_1(q)} |\psi\rangle_{\text{in}}$

Shear  $D_{f(\hat{q})} = e^{is\hat{q}^2}$ :  $e^{-is\hat{q}^2} \hat{p} e^{is\hat{q}^2} = \hat{p} + s\hat{q} = g(\hat{q} \sin \theta + \hat{p} \cos \theta)$

$$\begin{pmatrix} \hat{q}' \\ \hat{p}' \end{pmatrix} = \begin{pmatrix} \hat{q} \\ \hat{p} + s\hat{q} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ s & 1 \end{pmatrix} \begin{pmatrix} \hat{q} \\ \hat{p} \end{pmatrix} \quad \text{with } g = \sqrt{1 + s^2} \text{ and } \theta = \arctan s$$

Fourier transform:

$$\begin{pmatrix} \hat{q}' \\ \hat{p}' \end{pmatrix} = \begin{pmatrix} -\hat{p} \\ \hat{q} \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \hat{q} \\ \hat{p} \end{pmatrix}$$

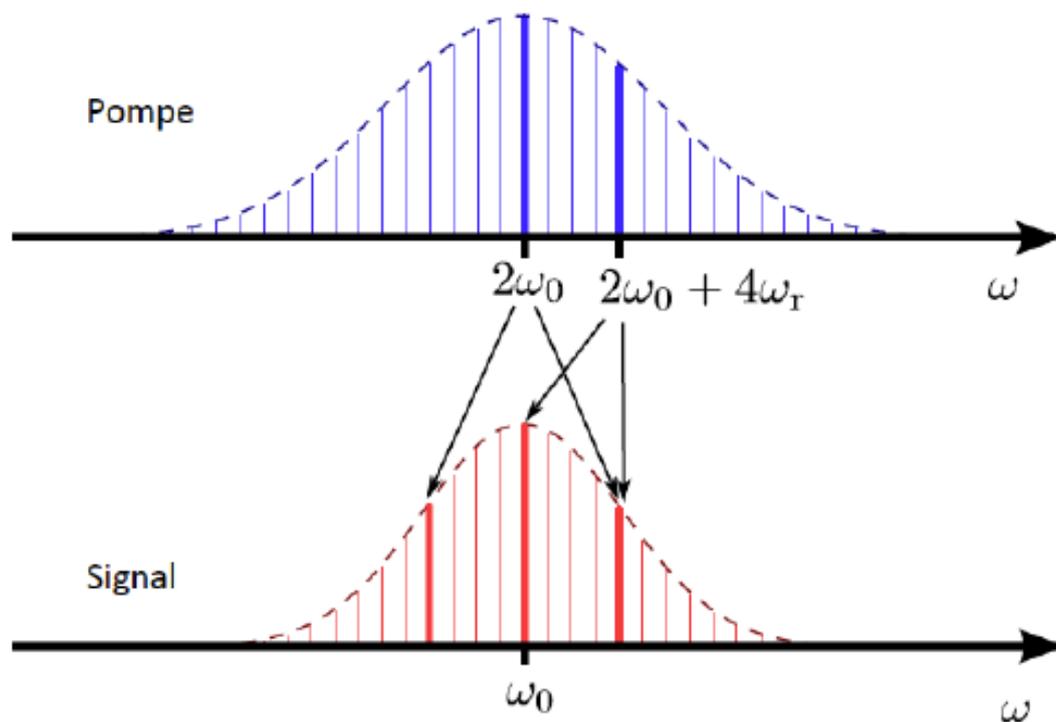
Hence, each step of the computation:

$$M(s) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ s & 1 \end{pmatrix} = \begin{pmatrix} -s & -1 \\ 1 & 0 \end{pmatrix} \quad \text{R. Ukai et al, PRA 81, 032315 (2010)}$$

In order to have  $M(s_2) = F$  we have to set  $s_2 = 0$

$$\rightarrow \theta_2 = \arctan s_2 = 0$$

# La conversion paramétrique est intrinsèquement multimode



Les **fréquences signal** sont toutes couplées entre elles grâce à la **pompe multimode en fréquence**

$$\hat{H} \propto \sum_{m,n} L_{m,n} \hat{a}_m^\dagger \hat{a}_n^\dagger + \text{h.c.}$$

signal

Matrice de couplage

$$\hat{S}_k = \sum_i U_k^i \hat{a}_i$$

$$\hat{H} \propto \sum_k \Lambda_k \hat{S}_k^\dagger \hat{S}_k + \text{h.c.}$$

Hamiltonien diagonal

$L$	$\Lambda_k$	$\hat{S}_k$
matrice	valeurs propres	supermodes
	20	s

# Un **SPOPO** se comporte comme une assemblée d'OPOs agissant chacun sur un supermode

